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Multiple Scattering Corrections for the Associated-Particle Neutron Time-of-Flight Technique

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Multiple Scattering Corrections for the Associated-Particle Neutron Time-of-Flight Technique

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Abstract

The computer code, MAGGIE, for the calculation of multiple scattering and sample attenuation in neutron differential cross-section measurements, has been revised and corrected. The particular case of the scattering geometry required by the associated-particle time-of-flight is considered in detail.

Key words: associated-particle, cross-section, Monte Carlo, multiple scattering, neutron, time-of-flight

MULTIPLE SCATTERING CORRECTIONS FOR THE ASSOCIATED-PARTICLE
NEUTRON TIME-OF-FLIGHT TECHNIQUE

Allan C. B. Richardson

I. INTRODUCTION

Measurements of fast neutron elastic and inelastic differential cross sections have, for many years now, usually been done using one of two time-of-flight techniques. The first of these requires a pulsed source of neutrons, and energy separation of the various neutron groups is then achieved by time correlation of the scattered neutrons with the incident neutron pulse. This technique has the advantage of flexibility of incident neutron energy and intensity, and the disadvantages of a relatively high time-correlated background and a low duty cycle. The other, and less commonly used, technique utilizes the detection of a charged particle from the neutron source reaction to tag the incident neutrons in time and direction. Energy separation of the various scattered neutron groups is then achieved by time correlation with the incident neutrons. This "associated-particle technique" has the advantages of very small time-correlated background, inherent absolute determination of the incident correlated neutron flux, and high duty cycle; but it suffers from limitations on the available neutron intensity and energy. However, at energies and intensities where this technique is applicable, it is the method of choice, since it is capable of yielding results of high accuracy without the ambiguities introduced by the time-correlated backgrounds and the massive shielding required by pulsed source techniques. The source reactions eligible for this method are those involving very light nuclei, and therefore capable of producing a light (and thus energetic) stable recoil nucleus. The $T(d,n)^4\text{He}$ reaction, producing 14-15 MeV neutrons at 90° over a rather wide range of incident particle energies, is most commonly used; other reactions that have been employed are $D(d,n)^3\text{He}$ and $T(p,n)^3\text{He}$, both of which produce lower energy neutrons. We will confine the discussion here to the $T(d,n)^4\text{He}$ case, although the method described is more generally applicable.

The associated-particle technique has scattering sample requirements that are quite different from those for a pulsed source. Instead of a relatively uniform incident neutron flux across the sample, the correlated neutron beam is highly directional. The angular distribution about the neutron beam line is usually well approximated by

$$I = I_0 e^{-\left(\frac{\theta}{\theta_0}\right)^2},$$

with θ_0 typically only a few degrees [1]. This property can be very useful [2]. It provides a high degree of neutron collimation without the need for massive collimators, which, in the case of 14 MeV neutrons, can produce substantial degradation of the initially monoenergetic neutron

beam. However, it forces the use of scattering samples of uniform thickness so that the cylindrical or spherical samples commonly used to simplify multiple scattering corrections are immediately ruled out. Otherwise a detailed knowledge of the neutron beam shape and extremely accurate alignment of this beam with respect to the scattering sample is needed. This alignment problem is further complicated because the center of the neutron beam slowly moves back as the neutron producing target ages with use. In measuring angular distributions of scattered neutrons it is of course also desirable to reduce the amount of scattering sample not directly in the neutron beam, so as to minimize multiple scattering.

A scattering sample in the form of a truncated cone, axis lying along the neutron beam, best satisfies all of these requirements. The origin of this cone is taken sufficiently far behind the source to allow for finite spot size on the neutron producing target, to provide some flexibility in alignment, and also to make allowance for changes due to target aging during a run. In order to make best use of the available neutron intensity, one must also use the thickest sample possible. The limit is set by either the angular resolution required at 90° or the time resolution required. A typical geometry is shown in figure 1.

In either case the resulting samples are sufficiently thick to require a careful multiple scattering correction. None of the analytical techniques [3], useful at energies up to a few MeV, are adequate at 14 MeV, the energy most commonly used for measurements of this type, since at energies above 6-7 MeV the diffraction peaks in the elastic angular distributions become too numerous. The only method of sufficient generality is the Monte Carlo technique. A survey of existing Monte Carlo codes revealed none for this particular geometry, but it was immediately apparent that the code "MAGGIE," developed by Parker, Towle, Sams, et al. at Aldermaston [4,5] contained all of the other elements important to such a calculation. For example, this code easily accommodates the neutron source distribution specified above. In addition, MAGGIE calculates an energy spectrum at each angle, so that false peaks due to double scattering can easily be identified. This sophistication is often useful, for example: the measurement of inelastic scattering from the very weakly excited $7.66 \text{ MeV } 0^+$ level in ^{12}C is easily confused by double scattering from the more easily excited $4.43 \text{ MeV } 2^+$ level in this nucleus.

A corollary need of any Monte Carlo neutronic calculation is an easily accessible, but at the same time sufficiently general, file program for the nuclear data required. This was available in a companion code to MAGGIE, entitled MOULD [6]. We have therefore modified the code MAGGIE so that it now handles the geometry required for associated-particle time-of-flight measurements. During the course of this modification we also corrected a few coding errors found in the original version, and made several additional modifications, principally updating the code to current computer syntax and capabilities.

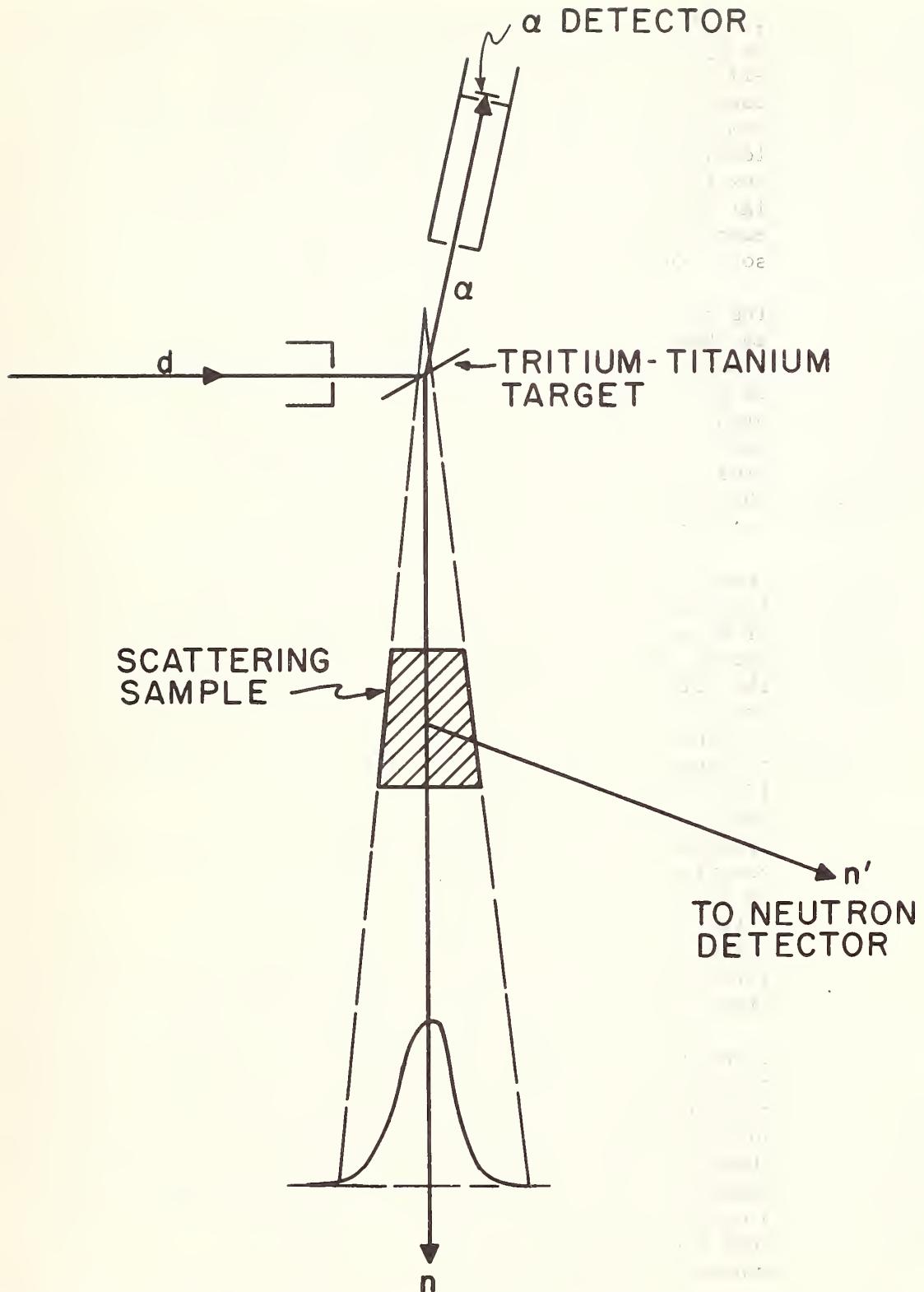


Fig. 1. The experimental geometry for associated-particle measurements at 14 MeV. The associated neutron intensity profile is indicated at the bottom of the figure.

Although a complete description of the code would be out of place here (see references 4 and 5 for details), a general overview of the program will facilitate later discussion of specific details. The point of view adopted is first to sample the scattering sample geometry and the available nuclear interactions using straightforward Monte Carlo techniques. Then, at each collision point, the weighted probability of scattering and escape to each of 33 detector positions (angles) in a half-plane lying to one side of the neutron beam axis is computed using the experimental data for elastic events and each of those inelastic processes of interest. The use of weighted probabilities of scattering to each of the detector positions at the final collision in the sample in place of a completely Monte Carlo approach results in a greatly reduced computation time. This is because the small solid angle subtended by each of the detectors (typically 10^{-3}) makes a final Monte Carlo scattering particularly inefficient. The output angular distributions obtained from applying this procedure to, typically, 1000 interacting neutrons are then reflected about the experimental input data and used as input for a second iteration. Two or three iterations are usually sufficient to obtain a convergent result.

In section II we give a description of the changes made in the code. First, the new coding for the scattering geometry used in the associated particle technique is described. Next we discuss a number of coding errors in the original version. Finally some changes are described that simplify the code and adopt it to FORTRAN V syntax. In section III we give some results obtained using data from the scattering of 14 MeV neutrons on natural carbon. A listing of those subroutines entirely rewritten or having extensive changes is given in Appendix A. In figure 2 the calling sequence for all of the components of the program is shown. A brief description of all of these subroutines appears in Appendix B.

II. DESCRIPTION OF MODIFICATIONS TO THE MONTE CARLO CODE MAGGIE

A. Changes Due to the New Scattering Sample Shape

Four subroutines are affected by changing the shape of the scattering sample. These are: 1) subroutine INPUT - those sections where scatterer parameters are read in and the flux attenuation factor is calculated are affected. The flux attenuation factor is defined as the ratio of incident flux along the axis of the sample to average flux in the entire sample. 2) CRNEU, the subroutine that creates incident neutrons at the entrance face of the scatterer by random sampling of the incident neutron spatial distribution. 3) TRACK, the tracking subroutine. 4) FPATH, the subroutine for calculating the probability of neutron escape from the sample in the direction of each of the assumed detector positions, for each collision point arrived at in TRACK.

It is worth noting here, although not necessary to what follows, that the data used for a) the neutron track lengths in the Monte Carlo sampling of the sample shape from subroutine EGMV, and b) the Monte Carlo sampling of reaction type and angular distribution at each collision by subroutine

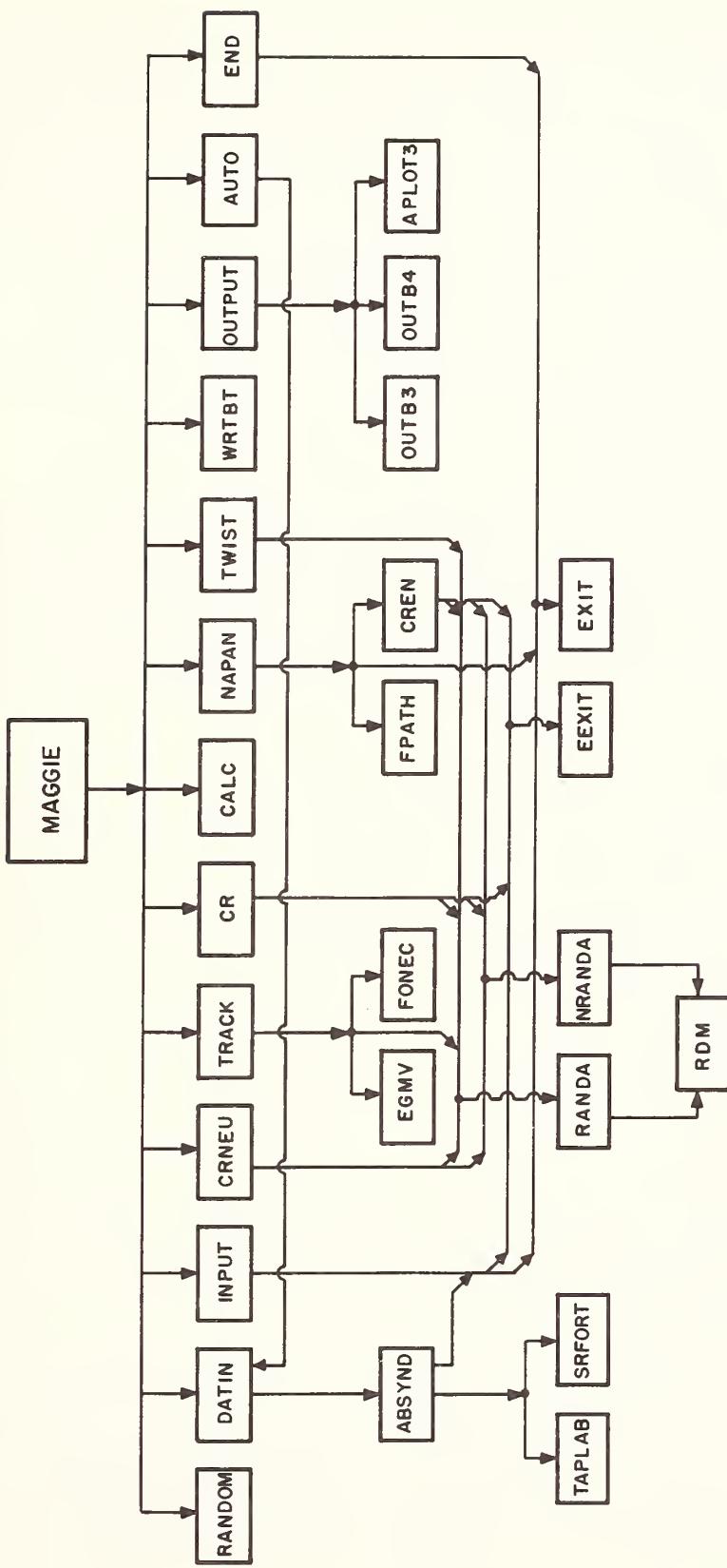


Fig. 2. The calling hierarchy for the present version of code MAGGIE.

CR, as well as c) the transmission probabilities in the direction of the detectors used in subroutines FPATH and NAPAN, are all drawn from the MOULD nuclear data file tape and not from input experimental data. The only experimental data employed are angular distributions used for the calculation of the scores for the relative probability of scattering to the various assumed detector positions in subroutine NAPAN. The elastic angular distribution (for the incident neutron energy only) used here is also automatically updated after each iteration by the subroutine AUTO. Thus it may be necessary to update the angular distributions on the MOULD file tape before and during the course of data-processing.

We now describe the changes due to the new scattering geometry in some detail for each of the subroutines affected.

1. INPUT

a. Cards 0207-0226 are changed to eliminate parameters required for samples consisting of concentric cylindrical shells, with common axes perpendicular to the beam axis, and to substitute those needed to characterize a truncated cone sample, axis lying along the neutron beam, and origin behind the neutron source. The new variables are HITE, FRAD, and ANGLE; the height, entrance face radius, and half-angle of the sample, respectively. HITE is immediately redefined as HITE/2, a more convenient quantity in subsequent calculations.

b. Card 0245, the calculation of the maximum angle subtended by the scattering sample at the neutron source, is changed to conform to the new geometry.

c. Cards 0301-0303 have been replaced. The correct expression for the flux attenuation factor for the new geometry is

$$\frac{F_o}{F} = \frac{K(o)}{\bar{K}} \cdot \frac{C}{N\lambda} \cdot \frac{h(r_1^2 + r_1 \Delta r + \Delta r^2/3)}{(1 - \cos \theta_m) d^2},$$

where h is the half-length of the scattering sample, r_1 the radius of its entrance face, Δr the difference in radii of the entrance and exit faces, and d the distance from the source to the entrance face of the sample. The remaining symbols are defined as in the original.

d. The original version of MAGGIE utilized a sample which was not symmetric about the axis of the incident neutron flux. Thus the experimental sampling of scattered neutrons in the detector plane was not truly representative of the scattering into 4π , which is employed in the program to calculate the flux attenuation factors and which are used in turn to infer the integrated cross-sections. This necessitated a small correction which was calculated with the help of a classification of the outgoing

Monte Carlo tracks vs. energy and the angle with respect to the scattering sample axis. Since the present scattering sample is symmetric about the incident neutron flux the experimental sampling of scattered neutrons in the detector plane is representative, and no correction is required. Accordingly, cards 0331-0338 in INPUT, cards 1544-1548 in DATIN, and cards 1735-1743 in MAGGIE are deleted. Card 1734 of MAGGIE is replaced by the statement

20 CONTINUE

The output of this table, by cards 3695-3715 of subroutine OUTB3, is also deleted.

2. CRNEU

Cards 3117-3119, 3140-3141, 3143-3147 and 3151 are replaced as shown in the listing. The new coding creates neutrons randomly scattered over the face of the conical scattering sample, in accordance with the specified spatial distribution. The old reference to multiple materials is not deleted, but the sample is now designated "material one." Similarly, the register containing neutrons which miss the sample is retained, although for this geometry misses occur with very low frequency.

3. TRACK

The entire subroutine has been replaced. The geometry is illustrated in figure 3. The subroutine is entered with the starting point (x_0, y_0, z_0) and the direction cosines $(\cos\theta_x, \cos\theta_y, \cos\theta_z)$ already defined. A random sampling of the neutron mean free path establishes the track length to the next possible collision. The track vector is then extended until it intersects the plane determined by the endface of the sample in the direction of travel of the neutron, and the path length from the starting point to this intersection is computed. A comparison of this intersection point with the radius of the endface establishes whether the track is in the direction of the endface or the curved surface of the sample, and the program branches accordingly. If the track passes through the endface, the path length calculated above is compared with the Monte Carlo track length to the next collision and the coordinates of collision or escape from the surface of the sample as well as the time elapsed along the track are computed in a straight-forward way.

We consider now the procedure used for a neutron headed toward a curved surface of the scattering sample. The coordinates of the endpoint of the previously determined Monte Carlo track length are first found. These are used to calculate the projected distance, perpendicular to the symmetry axis, of this endpoint from the symmetry axis, as well as the length of the radius of the sample lying along this projection. Comparison of these quantities determines whether a collision occurs within the sample or not. If so, the coordinates are already known and the time elapsed along the track is all that remains to be calculated. If not, it is necessary to calculate the coordinates of the neutron's exit from

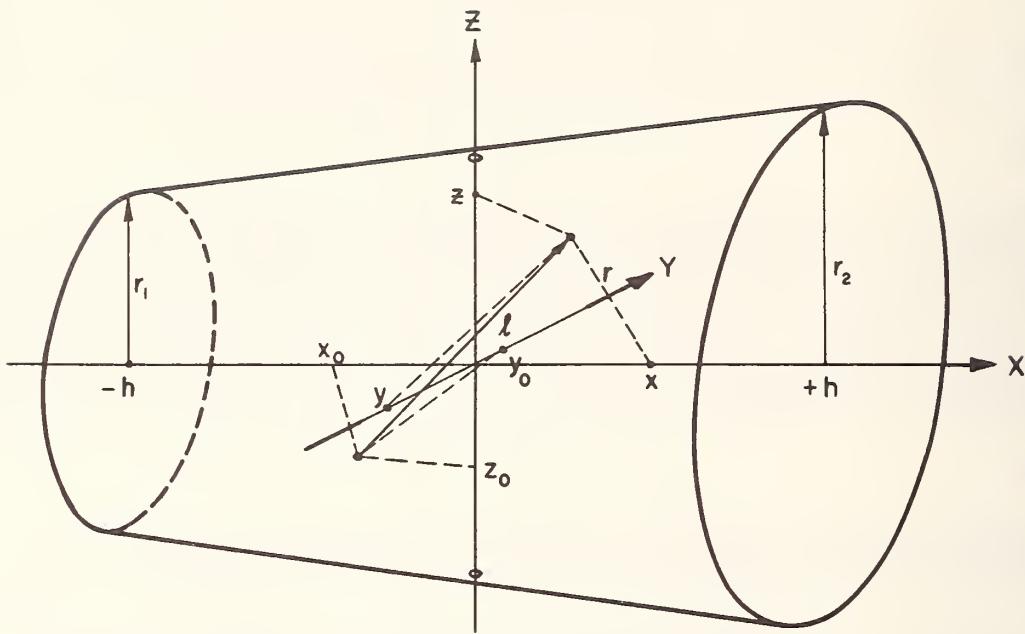


Fig. 3. The geometry for subroutine TRACK. The trajectory shown is that of a neutron which, after collision at point x_0 , y_0 , z_0 escapes from the curved surface of the sample at point x , y , z .

the surface of the cone. Designate ℓ the path length to the surface and x, y, z the coordinates of the intersection of the track with the surface. We have, referring to figure 3,

$$x - x_o = \ell \cos \theta_x$$

$$y - y_o = \ell \cos \theta_y$$

$$z - z_o = \ell \cos \theta_z ,$$

so that the projected distance, r , from the symmetry axis of the cone to the point (x, y, z) is given by

$$\begin{aligned} r^2 &= y^2 + z^2 \\ &= \ell^2 \sin^2 \theta_x + 2\ell(y_o \cos \theta_y + z_o \cos \theta_z) + y_o^2 + z_o^2 , \end{aligned} \quad (1)$$

where we have used the identity

$$\cos^2 \theta_x + \cos^2 \theta_y + \cos^2 \theta_z = 1 .$$

Since we know the parameters of the cone we can also calculate r from the x coordinate (again see figure 3.):

$$r = r_1 + (h+x) \tan \theta_o , \quad (2)$$

where θ_o is the half-angle of the cone. Eliminating r from eq^{ns.} (1) and (2) and arranging the terms as a quadratic in ℓ , we obtain

$$\begin{aligned} &\ell^2 \left[1 - \cos^2 \theta_x (1 + \tan^2 \theta_o) \right] \\ &+ 2\ell \left[y_o \cos \theta_y + z_o \cos \theta_z - \cos \theta_x (r_1 \tan \theta_o + h + x_o) \tan^2 \theta_o \right] \\ &+ \left[y_o^2 + z_o^2 - (r_1 + (h+x_o) \tan \theta_o)^2 \right] = 0 . \end{aligned} \quad (3)$$

The positive root of this equation is the required path length to the surface, and the subroutine FONEC is then called to provide the x, y , and

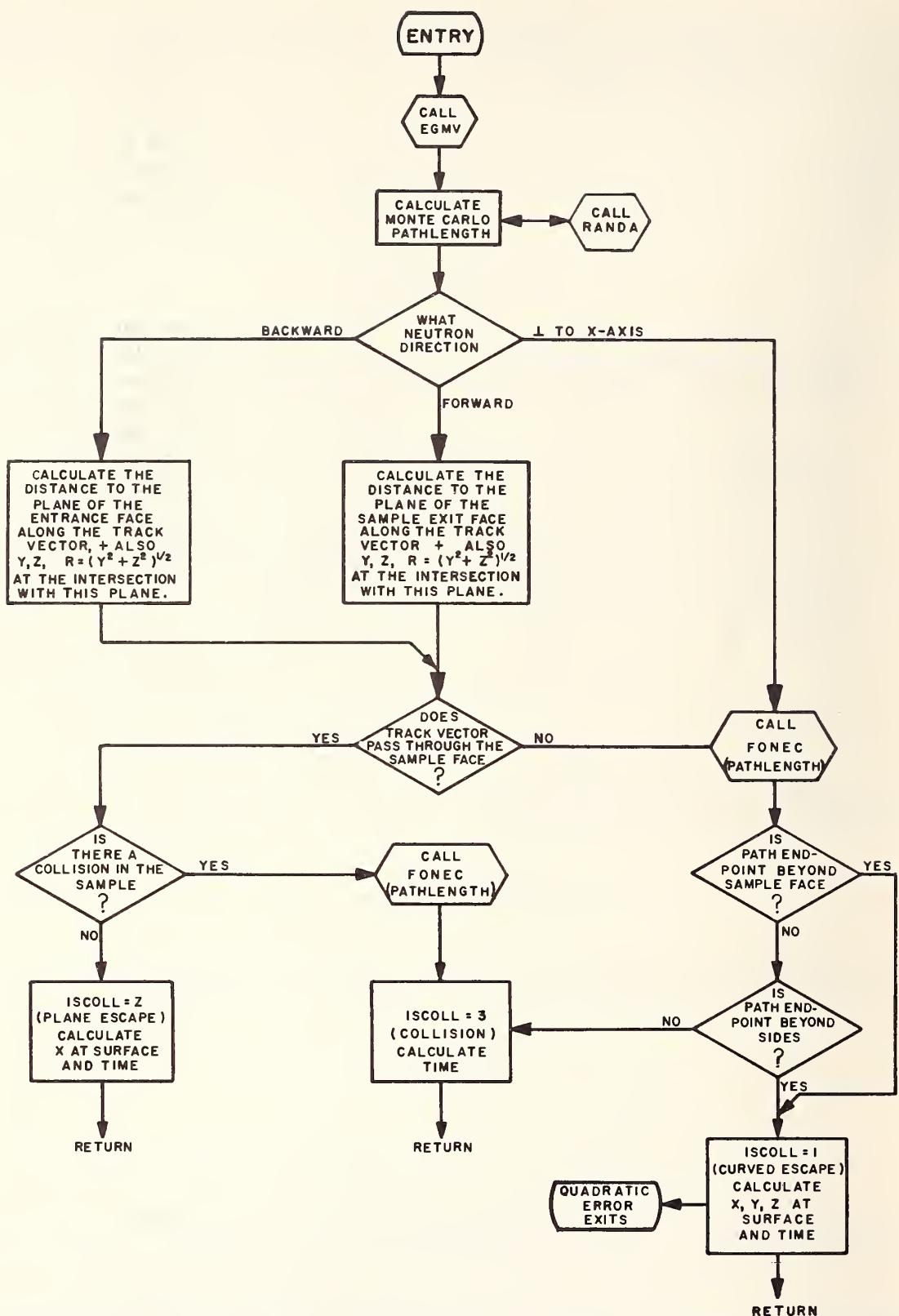


Fig. 4. Subroutine TRACK. The subroutines EGMV and FONEC return the neutron mean free path and the coordinates at the end of a track, respectively. RANDA returns a random number.

z coordinates. Error prints are provided for the cases of two positive, two negative, imaginary, or indeterminate roots as solutions to eq. (3). However none of these situations has occurred in many tens of thousands of tracks, although at least some of them are mathematically possible. Once ℓ has been calculated, determination of the coordinates of escape and the track time are straightforward. The subroutine also returns an index that indicates the fate of the neutron, i.e., curved escape, plane escape, or collision. A block diagram is shown on figure 4.

4. FPATH

This subroutine has also been entirely replaced. The same simplifying assumption is made as in the original coding--that is, the dimensions of the scattering sample are considered to be negligible compared to the flight path to the detector. This has the effect that the angle and flight path to any particular detector position may be considered to be the same from any scattering point in the sample. The quantities to be calculated are the neutron path length, p , in the sample from any point (x, y, z) to the surface of the sample and the scattering angle, θ_d , in the direction of a particular detector, at the angle ψ . The geometry is shown in figures 5 and 6. All paths are assumed to be parallel to the x - y plane, and all detector positions satisfy the condition $y \geq 0$.

We note that the locus of the intersection of a cone with a plane parallel to its axis is an hyperbola, and write, using the notation shown in figure 5,

$$\frac{x^2}{a^2} - \frac{y^2}{b^2} = 1 \quad ,$$

$$\text{or} \quad y^2 = x^2 \tan^2 \theta - z \quad , \quad (4)$$

$$\text{since} \quad b = z \quad \text{and} \quad a = z/\tan\theta.$$

Now, referring to figure 6, we may also write

$$Y = p \sin\psi + y$$

$$\text{and} \quad X = p \cos\psi + x' \quad , \quad (5)$$

$$\text{where} \quad x' = R + h + x = \frac{r_1}{\tan\theta} + h + x.$$

Again we reduce the problem to a quadratic in path length, p , after

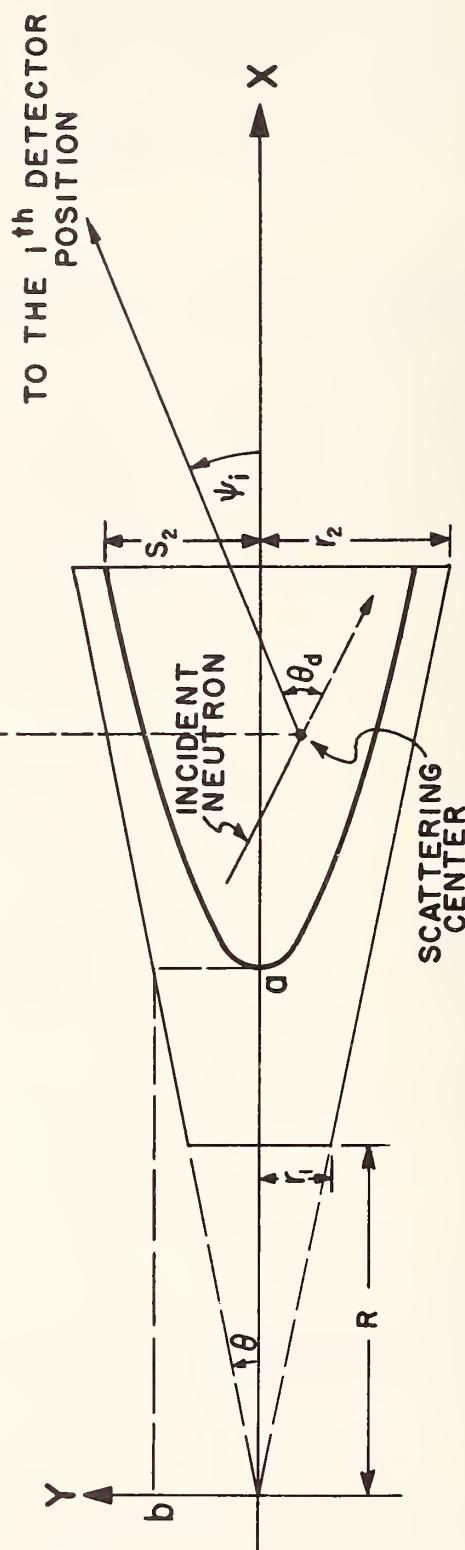
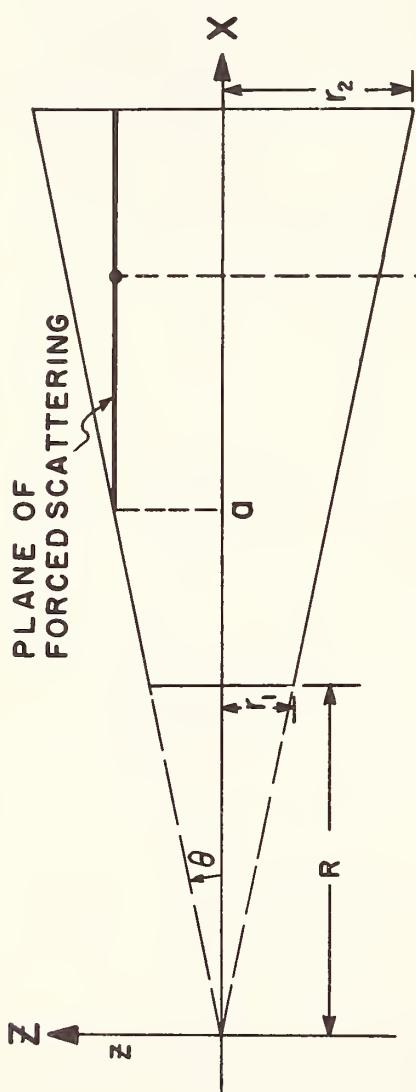
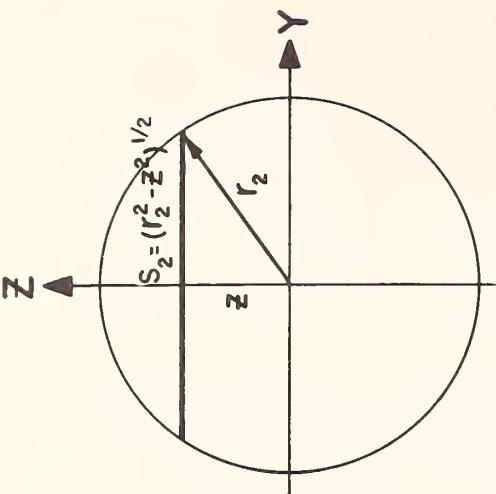


Fig. 5. The geometry for subroutine EPATH. The plane of forced scattering is located at a height $r_1 < z < r_2$ in this instance. The track of the incident neutron does not, in general, lie in this plane as shown.

TO THE i^{th} DETECTOR POSITION

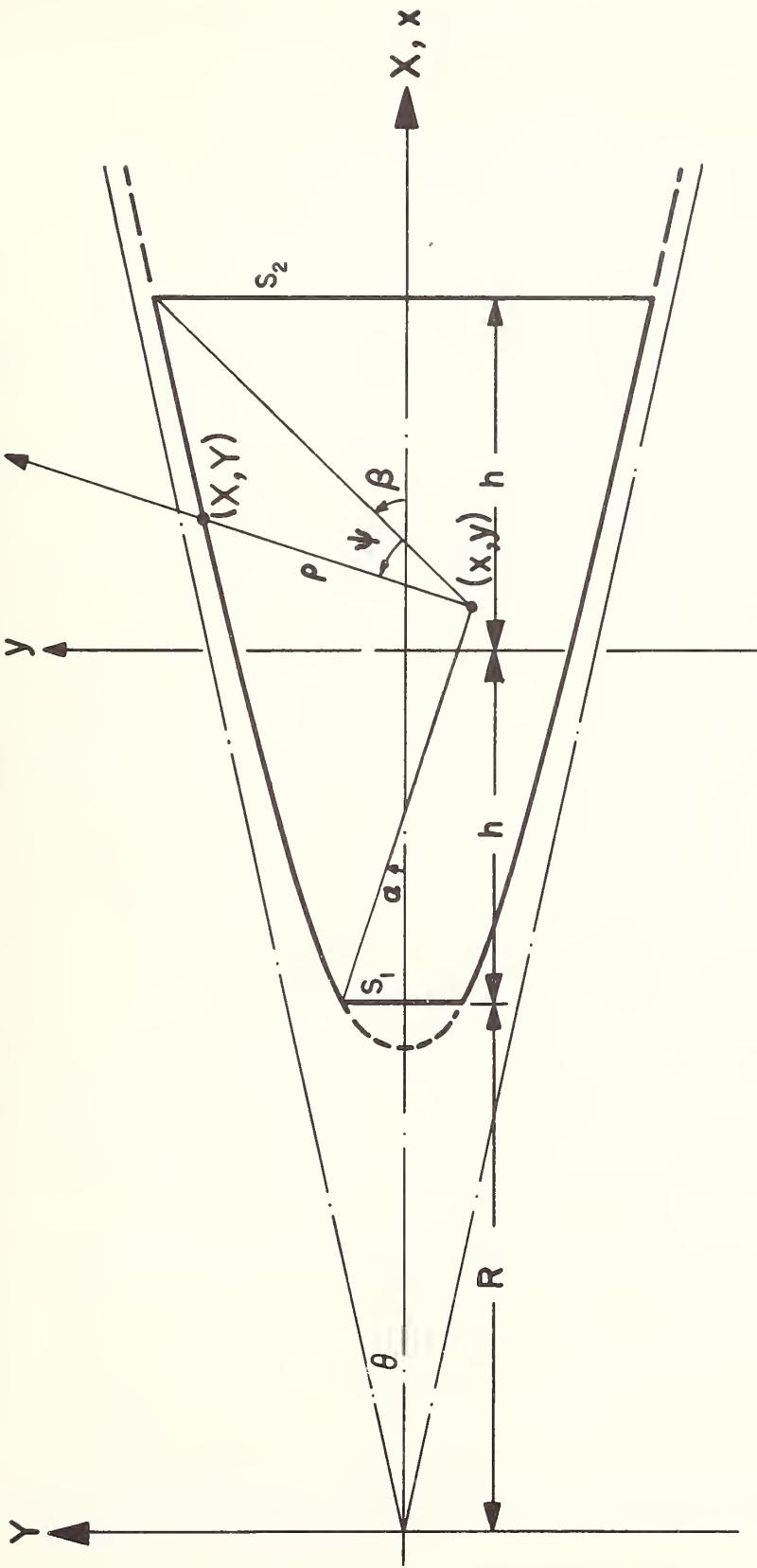


Fig. 6. A section through the scattering sample at height $z < r_1$. The upper case quantities refer to the axes X and Y , whose origin is that of hyperbola defined by the intersection of the plane (x, y) with the surface of the scattering sample.

combining eq's (4) and (5) to eliminate the variables X and Y to obtain:

$$\begin{aligned} & p^2 (\sin^2 \psi - \cos^2 \psi \tan^2 \theta) \\ & + 2p(y \sin \psi - x' \cos \psi \tan^2 \theta) \\ & + (y^2 - x'^2 \tan^2 \theta + z^2) = 0 . \end{aligned} \quad (6)$$

For each value of z there will be a pair of angles, α and β (see figure 6), which define the limits of the hyperbolic curve. These angles are given by

$$\begin{aligned} \tan \alpha &= \frac{(r_1^2 - z^2)^{1/2} - y}{h + x} \\ \tan \beta &= \frac{(r_2^2 - z^2)^{1/2} - y}{h - x} . \end{aligned} \quad (7)$$

The subroutine is diagrammed in figure 7. After testing for the special case of exit perpendicular to the beam axis, the program tests, using eq's (7), for exit through the endfaces vs. the curved sides for the exit angle, ψ , and branches accordingly. The path length in the former case is a straightforward calculation, and for curved escape is the positive root of eq. (6). It should be noted that in the limiting case of cylinder ($\tan \theta = 0$) this subroutine, unlike all of the others, does not work. A modified subroutine for a cylindrical sample is given in the Appendix following the listing of FPATH for a truncated cone.

B. Correction of Coding Errors [7]

1. INPUT

a. Statements resulting in $LGR=0$ in the table look-up for the mean free path (following card 0308) have all been changed to give $LGR=1$. $LGR=0$ references not a mean free path, but instead the last tabular value of collision probability in the previous MOULD table. The revised look-up supplies the mean-free path at the lowest tabulated energy for any neutron at or below that energy.

b. On the card following statement 5001, LRG has been replaced by the correct variable, LGR.

c. Card 0441 is replaced by

IMINM = IMAXM+1 .

This change yields sequential storage of supplementary ranges of experimental angular distributions, and thus avoids loss of needed storage space in the array SUPVAL (I,J).

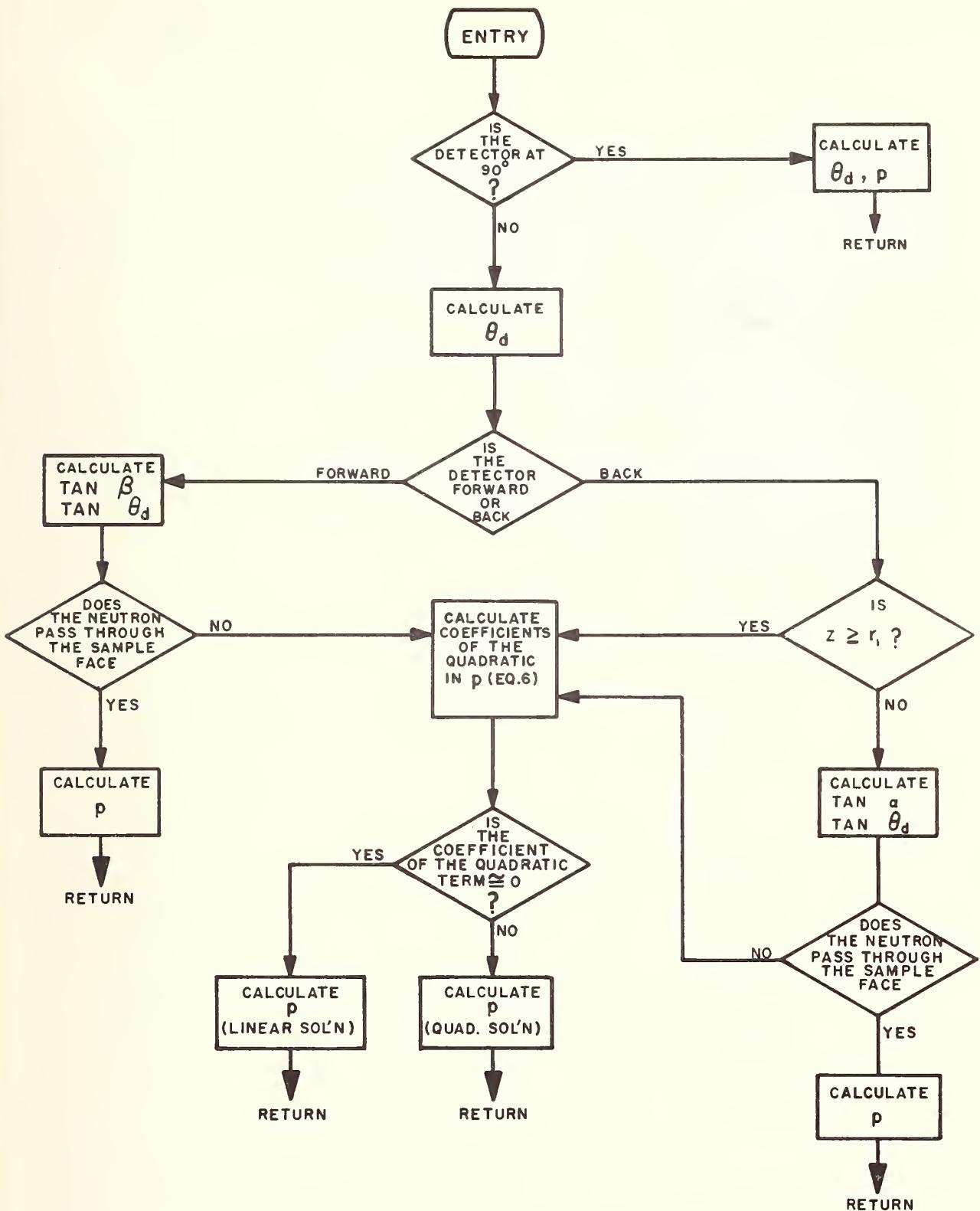


Fig. 7. Subroutine FPATH, for a truncated cone. The subroutine for a cylinder is similar.

2. EGMV

LGROUP=0 has been eliminated for the reason given in 1a. above.

3. NAPAN

a. The angle used for the interpolation of the experimental distributions, as well as for calculation of the neutron energy after collision is the laboratory angle, calculated in the subroutine FPATH. For this reason the experimental angular distributions must be read in at laboratory angles. Also the center-of-mass energy calculation on cards 3494 and 3497-3501 (in the case of center-of-mass MOULD angular distribution tables) is in error. This has been changed to a calculation in laboratory coordinates by replacing card 3494 by

GO TO (36,11,11,11,11,11,11,11,36,11,11,11,11,11,11) NFORM

and deleting cards 3497-3501.

b. The sequence of cards 3502-3512, which picks the appropriate experimental angular distribution values, does not refer to the correct angular distributions, nor to the correct neutron energy--that before collision. It has been replaced by the sequence:

```
11      IF(IMAX(J).LE.0) GO TO 15
        IMINM=IMIN(J)
        IMAXM=IMINM+IMAX(J)-1
        DO 50 JJ=IMINM,IMAXM
        IF(EIN-ENVAL(J,JJ)>50,51,52
          50      CONTINUE
          51      KK=JJ
          GO TO 19
          52      IF(JJ.LE.IMINM) GO TO 15
                    KK=JJ-1
```

c. On the card following statement 5001, LRG has been replaced by the correct variable, LGR.

d. LGR=0 has been eliminated for the reason given in 1a. above.

e. The option that provides for suppression of printout of various scores, through setting the indicators IANALA, IANALB, IANALC, and IANALD negative, should not suppress storage of these scores by NAPAN. The time saved is negligible and the elastic scores, DEB(I,L), are required by the automatic iteration scheme, AUTO. Cards 3528-43 have been revised to eliminate the dependence on these printout indicators.

4. ABSYND

LGROUP=0 has been eliminated for the reason given in 1a. above.

5. CREN

For the same reason set out in section 3a. above, the center-of-mass energy calculation (cards 2960-66) has been replaced by a calculation in laboratory coordinates by deleting cards 2961-66 and replacing card 2960 with

880 GO TO 800 .

6. OUTPUT, OUTB3, OUTB4

The sum over energy, at each detector angle, stored in the vector DES(J,33) is required by subroutine AUTO. So as to allow calculation of this quantity in subroutine OUTB4, even when no B4 printout is requested, a print indicator, IFPR, has been added to the call for subroutine OUTB4. This is accomplished by the changes shown in the listing for cards 3721, 3808A, 3873, and 3975-4002.

7. AUTO

a. The quantity "chi-squared," on card 2084, should be

$$\chi^2 = \sum_{1}^{33} \frac{\left(\sigma(\theta) \exp^{-J(\theta)} \text{calc} \right)^2}{\sigma(\theta) \exp} ,$$

and not as formerly written, with the denominator squared.

b. It should be noted that the normalization of the Monte Carlo output on cards 2065-2068 and 2088 is not, in general, correct; but only holds for detectors positioned at equal increments of $\cos\theta$ between -1 and +1. For this reason, if subroutine AUTO is used, the input detector positions must satisfy this criterion.

8. FPATH

The original version of this subroutine contained two errors. These do not apply to the new version of FPATH reported here, but are listed for completeness. Card 3356 of the original should read

SINOM = SQRT(1-COSOM**2) ,

card 3358 be deleted, and a new card inserted following card 3359:

FP1 = WORK16**2+WORK17**2-FP2**2 .

C. Other Modifications

1. General

a. Common storage assignments have been handled by combining all common statements into one package, processed by the UNIVAC 1108 "Procedure Definition Processor" using the assembler directive FCOPY (Fortran copy). This package, designated CINC1, is then included at the time of assembly in all subroutines making use of common variables by means of the statement

```
INCLUDE CINC1
```

placed immediately following the subroutine name definition. This process, which effects a considerable economy in the size of the source deck and listing, is characteristic of the UNIVAC 1108 Assembler and FORTRAN V; however, equivalent procedures are often available to other systems. In some subroutines dimension statements still appear for those few variables not in common storage. A block of common storage has been specifically assigned to the subroutine ABSYND variables DATA and IDATA. This carries the dummy label BLANK. The details of these changes in memory storage allotment will be obvious upon examination of the listings.

b. The size of many modern computers obviates the need for linkage. We have deleted references to CHAIN(I,J) and incorporated the balance of subroutine PRELUDE into the main program MAGGIE. Similarly, references to CHAIN(I,J) in subroutines DATIN and AUTO have been replaced by the appropriate CALL and/or RETURN statements.

c. The final F in the names of all library functions (such as SIN, COS, EXP, MIN) has been deleted, as it is not compatible with FORTRAN V. In addition, such functions as FLOAT and INT have been eliminated by using the mixed expressions allowed by FORTRAN V.

d. Disc storage read and write statements have been modified in the main program MAGGIE and subroutine WRTBT to conform to 1108 FORTRAN V usage.

e. A number of formats have been altered. Some of the changes are merely different spacing options, but others are required by the new scattering geometry. A list of card numbers of the affected statements follows:

0226	0226A	0226B
0249	0258	0345
0410	0450	0451
1719	1963	1972
2111	2115	3638
3639	3828	3833

On card 2735 in subroutine CR the variable N has been replaced by the correct variable NSECS.

f. A new random number generator, subroutine RANDOM, entry RDM, has been incorporated, and the required changes in the functions RANDA and NRANDA are shown in the listings.

2. MAGGIE

a. All of the subroutine PRELUDE, except for calls to subroutines CLOCK and CHAIN, has been inserted following card 1654. A call to subroutine RANDOM immediately following the read-in of OCT, the starting value of the random number generator, initializes this generator. OCT is printed out following the call to subroutine INPUT.

b. All references to subroutine CLOCK have been deleted.

c. Card 0530 is eliminated by changing cards 1776, 1778, 1787, 1789, and 1792 to read

GO TO 1 .

3. ABSYND

a. Card 0525, which assigns the logical tape unit carrying the MOULD nuclear data tape, now reads

NUCDAT=9 .

On cards 0525, 0530, and 1394 the variable name TAPE has been changed to ATAPE to avoid confusing the compiler. These changes may not be required by other installations.

b. Cards 0553 and 0554 set all storage for the variables DATA and IDATA to zero before each ABSYND run.

c. Card 1398 is modified so as to print only those action numbers processed.

4. CR

The variable names NACT and Q have been changed to NACTV and QV, so as to avoid conflict with the array names NACT(I) and Q(I). Similarly, the variable name COS is changed to COZ, so as to make the library function COS available.

5. CREN

a. The variable names NACT and Q have been changed to NACTV and QV, for the reason given above.

b. Cards 2832-2837, 2878-2881, and 2911-2915 have been deleted, since they are not needed. In order to accomodate these deletions card 2910

has been amended to read

IF(NVCOS.LT.1) GO TO 208

,

and card 2916 to read

P = (AS+SQRT(AS**2+A*QV*(1+A)/EIN+A**2-1))/(1+A) .

6. AUTO

a. The multiple elastic sums are not required, and so cards 2069-2072 are deleted.

b. The quantity RATIO is not used, and has been eliminated from cards 2077 and 2078.

c. The do-loop 2083-2085 is redundant, since ELM2(J) is also set at card 2097. Cards 2083-2085 are therefore deleted, and card 2080 is modified to reflect this change.

d. The three do-loops in the sequence 2087-2098 have been combined into a single loop.

e. The variable name EXP(J) has been changed to EZP(J) to avoid conflict with the library function EXP.

f. Often it is desirable to separate an iteration procedure into two or more consecutive runs. Cards 2073-2082 have been modified to allow a run to be made using the partially corrected output of a previous run as input, instead of the experimental data. When this option is to be used the second field of the input card specifying the number of iterations, NTERM, should be non-zero. This is then followed by the cards specifying the experimental values for the angular distribution, FCVAL, which are removed from their usual position and replaced by the partially corrected output of the previous run.

g. In the case of rapidly varying angular distributions such as at 14.1 MeV, the usual iteration procedure is not as rapidly convergent as the "physical" method. For this method the calculated multiple scattering is first subtracted from the experimental input, and then the balance of the iteration (sample attenuation) is performed as usual by reflecting the output for single scattered neutrons about the input. The new coding on cards 2091-2097 reflects this change. For the first two iterations the iteration improves only the multiple scattering, the correction being applied both times to the experimental input. After two iterations multiple scattering is well enough known to allow the iterative procedure full play, so the correction is applied each time to the previous input, instead of the experimental values. This procedure yields much more rapid convergence, three iterations providing better convergence than six of the previous method.

III. AN APPLICATION TO 14.1 MeV NEUTRON SCATTERING

The code has been applied to a typical associated-particle scattering geometry for 14.1 MeV neutrons on carbon. The scattering sample was a truncated cone of half-angle seven degrees, length 3.193 cm., and entrance face radius 2.059 cm., with its center located 20 cm from the neutron source. The experimental input used was simulated using published results for carbon at this energy. The results of a typical run are shown in Table 1, and on figures 8 and 10. The cross-section data used for the Monte Carlo scattering were those of Slaggie and Reynolds [8]. In Table 1 the quantity $\sigma(\theta)$ represents the true angular distribution, and the total and multiple outputs are the results calculated using it and the specified sample shape. The code varies $\sigma(\theta)$ until the total elastic output matches the experimental input to the accuracy required. The calculated angular distribution for neutrons undergoing multiple scattering is shown in figure 8. It can be seen that the result converges quite rapidly. Further iterations resulted only in statistical variations about the values obtained after three iterations; in fact, the largest change after two iterations, that at $\cos \theta = 0.625$, represents only a 1.5% change in the cross section.

In figure 9 we show the part of the correction, exclusive of multiple scattering, that depends upon sample shape. Also displayed for comparison is a calculation of the effect expected due to attenuation in the sample for an isotropic angular distribution. As can be seen, this correction is not independent of the input angular distribution. The quantity plotted, $\Delta\sigma/\sigma$, is the difference between the singly scattered output and the input, divided by the input, for each angle. This geometrical correction is contributed to about equally by varying path length in the sample and the change in the total neutron cross section with energy as a function of the neutron scattering angle. The correction varies from a few percent to about ten percent, and is directly dependent upon the accuracy with which the total neutron cross section is known over the range of energy exhibited by the neutron recoil. On the other hand, the correction for multiple scattering is larger, ranging from about 50% at the backward minimum to 4% at zero degrees, but depends mainly upon the experimental input data. Over most of the practical angular range of measurement it is a quite appreciable 15-20%. Figure 10 shows the input distribution used and the corrected output distribution obtained after three iterations of approximately 2000 interacting neutrons each. The total correction is, of course, largest in regions where the cross section exhibits minima. It varies from a few percent to a maximum of only twenty percent at measurable angles, since the geometrical and multiple scattering corrections tend to be in opposite directions, except at the extreme backward angles. Several examples of spectra showing elastic-inelastic and multiple inelastic effects in the energy spectra are given in the original references [4, 5]; the present version of MAGGIE also yields similar results.

TABLE 1

The scores for a typical Monte Carlo run for carbon at 14.1 MeV.

COUNTER NUMBER	COS θ	$\sigma(\theta)$	TOTAL ELASTIC OUTPUT	MULTIPLE ELASTIC OUTPUT	EXPERIMENTAL INPUT
1	-1.0000	.023	.044	.021	.044
2	-.9375	.075	.092	.022	.093
3	-.8750	.093	.111	.024	.111
4	-.8125	.100	.120	.026	.121
5	-.7500	.110	.131	.028	.131
6	-.6875	.127	.148	.030	.148
7	-.6250	.152	.170	.031	.171
8	-.5625	.176	.190	.033	.191
9	-.5000	.202	.212	.034	.212
10	-.4375	.225	.231	.035	.232
11	-.3750	.240	.246	.036	.246
12	-.3125	.248	.253	.037	.254
13	-.2500	.245	.251	.037	.251
14	-.1875	.233	.240	.038	.241
15	-.1250	.216	.225	.038	.226
16	-.0625	.192	.205	.038	.205
17	.0000	.173	.183	.038	.183
18	.0625	.145	.155	.037	.155
19	.1250	.125	.135	.037	.135
20	.1875	.114	.125	.037	.125
21	.2500	.107	.123	.039	.122
22	.3125	.110	.127	.041	.126
23	.3750	.124	.142	.044	.140
24	.4375	.154	.170	.049	.167
25	.5000	.203	.215	.055	.212
26	.5625	.281	.285	.063	.281
27	.6250	.412	.404	.074	.400
28	.6875	.628	.605	.087	.600
29	.7500	.969	.929	.104	.925
30	.8125	1.486	1.435	.125	1.435
31	.8750	2.274	2.218	.151	2.219
32	.9375	3.451	3.380	.181	3.387
33	1.0000	5.045	5.197	.216	5.062

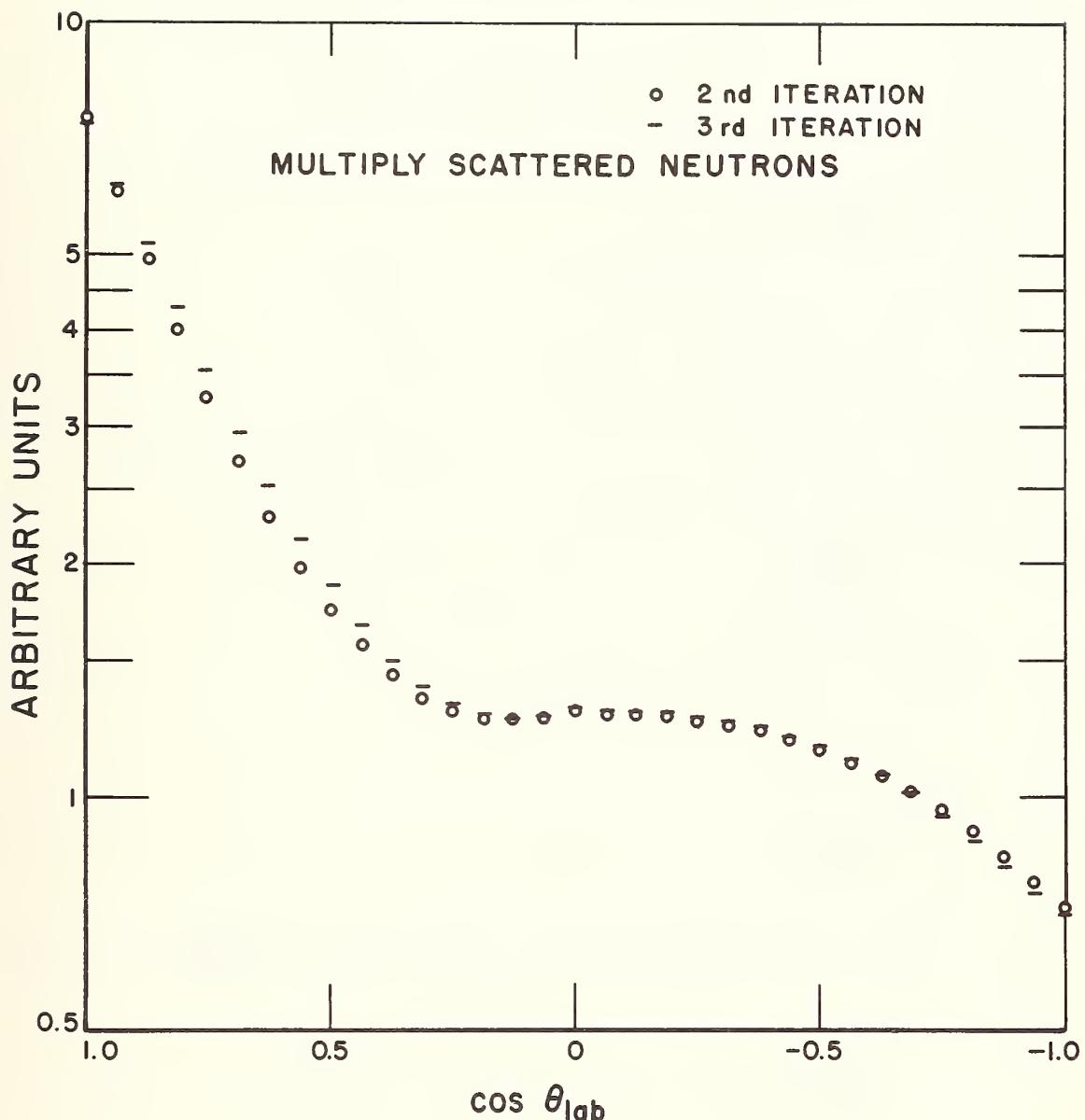


Fig. 8. The angular distribution of multiply scattered neutrons on carbon at 14.1 MeV for two successive iterations. The data for the third iteration are also given in Table 1, after suitable normalization.

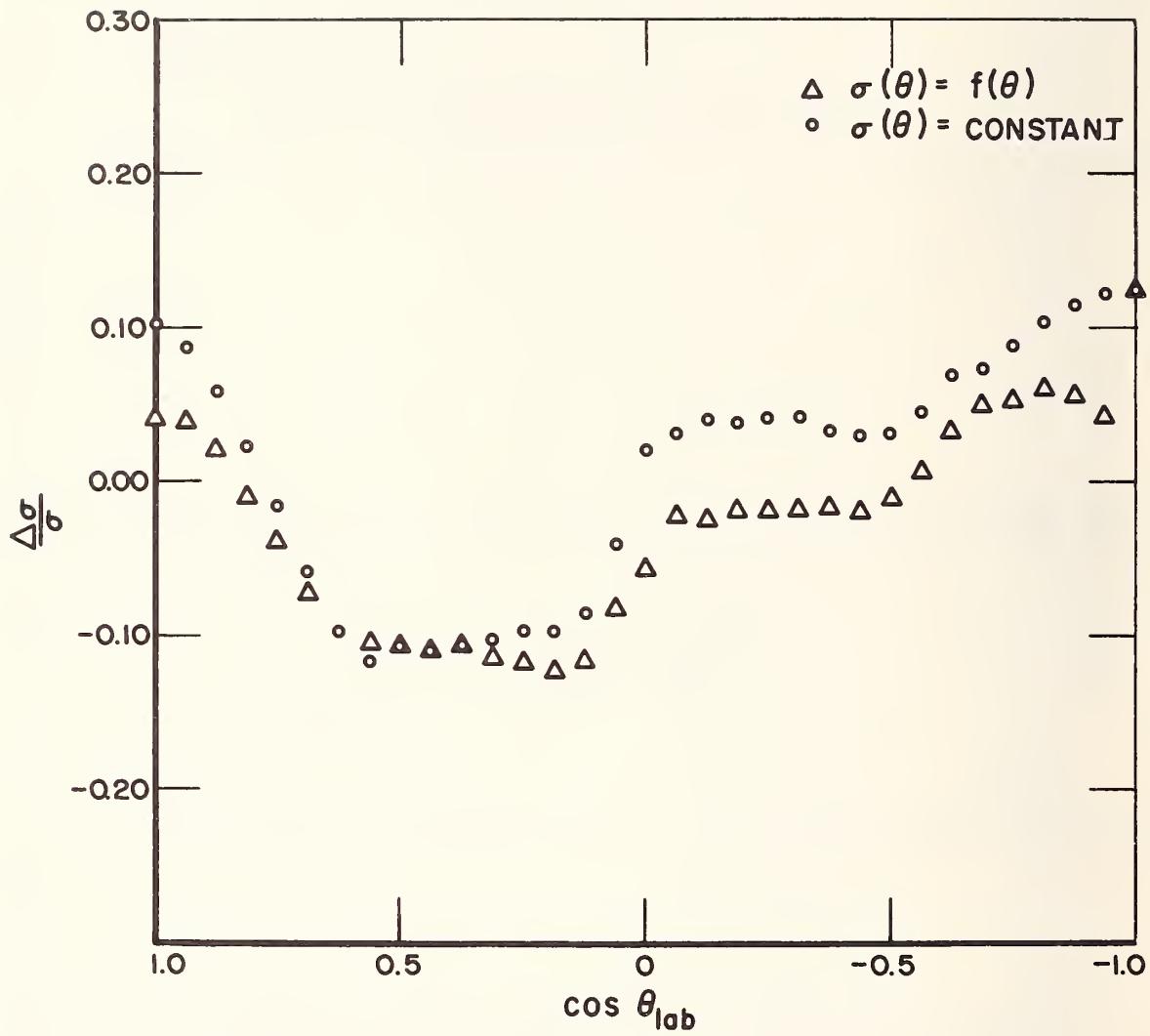


Fig. 9. The correction due to sample shape. The quantity $\Delta\sigma/\sigma$ is the change in the normalized angular distribution after subtracting multiple scattering.

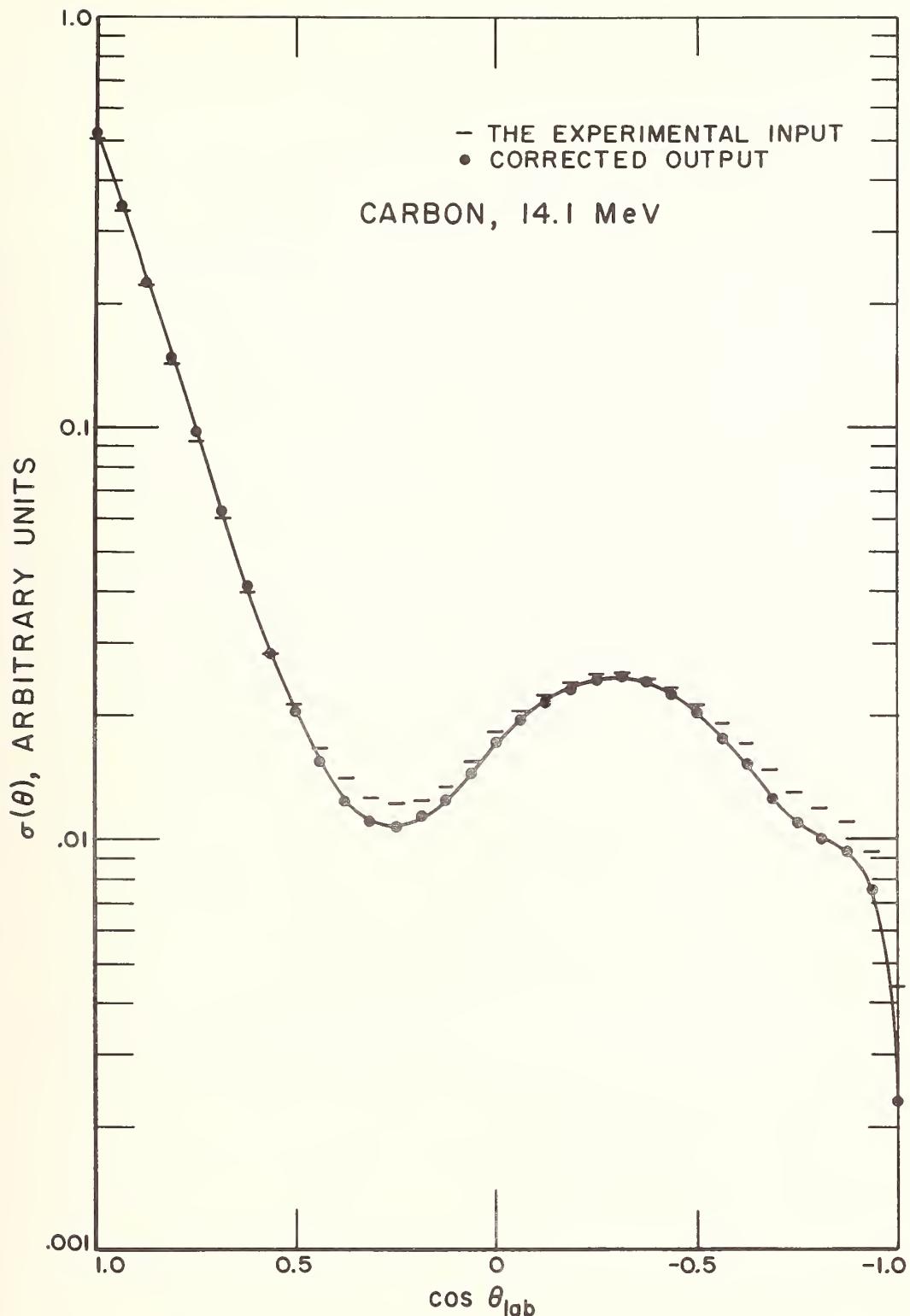


Fig. 10. The experimental input and corrected output for carbon at 14 MeV.

REFERENCES

1. Monier, L. F. C., Tripard, G. E., and White, B. L., Nucl. Instr. and Meth. 45, 282 (1966).
2. Marshak, H., Richardson, A. C. B., and Tamura, T., Phys. Rev. 150, 996 (1966).
3. A useful summary is given by M. Walt in Fast Neutron Physics, Vol. II, Marion, J. B., and Fowler, J. L., Ed. (1960).
4. Parker, J. B., Towle, J. H., Sams, D., and Jones, P. G., Nucl. Instr. and Meth. 14, 1 (1961).
5. Parker, J. B., et al, Nucl. Instr. and Meth. 30, 77 (1964).
6. Parker, K., AWRE report 0-70/63; Kerr, W. M. M., AWRE report 0-81/64 (1964); Miller, S. M., and Parker, K., AWRE report 0-55/65.
7. Some of these coding errors have also been previously noted in private communications. We are indebted to J. B. Parker, Aldermaston, for noting the substitution of REDEN for EIN in subroutine NAPAN, as well as the errors in subroutine FPATH. H. Horstmann and H. Schmid, Geel, have noted, in addition, the difficulty involving the center of mass calculation of neutron recoil energy in subroutines NAPAN and CREN.
8. Slaggie, E. L., and Reynolds, J. T., KAPL-3099, (1966).

APPENDIX A

Listings of new and extensively revised coding.

<u>Name</u>	<u>Page</u>
CINC1	28
MAGGIE 3A	30
INPUT	36
CRNEU	44
TRACK	46
NAPAN	49
FPATH (truncated cone)	53
FPATH (cylinder)	55
OUTPUT	56
AUTO	58
RANDOM	61
RANDA	61
NRANDA	61

A complete listing is available by request from the Center for Radiation Research, Neutron Physics Section, National Bureau of Standards, Washington, D. C. 20234.

```

CINCI* FCOPY
C
C COMMON STORAGE ALLOCATION.
C
C COMMON LOC,OCT
C DIMENSION DATA(1),IDATA(1)
C EQUIVALENCE (DATA, IDATA)
C COMMON IDATA,INIT,NEGS,UMAX,XLYMIN,IXMAT,IDICE,NMATS,EMC(257),
1 PATH,SPEED,MAT,EIN,UIN,LGROUP,ROOTE,WIN,WOUT,NSECS,
2 COSPHI(100),EOUT(100),UOUT(100),ISORAN(100),NOLAW(100),
3 NONUC(100),NOACT(100),NREACT(100),NORNGE(100),ALF,
4 ALFP(20),WMIN,WMAX
C
C MAGGIE NUCLEAR DATA AIDS.
C
C COMMON NUCLID,NAC,AS,ENER,NFORM,QV,ATOM1(24),Q(24),PART1(24),
1 IXACT1(24),IXNUC1,FNEGS,NUCS,NUCL(24),NACT(24),NAC1(24),
2 NANAL,NENSP,NOCOU,NOGR,LAWREV(50),LWNO(50),NLWREV
C
C BIRTH STORES AND REGISTERS.
C
C COMMON BIRTH(5500),IBIRTH(1500)
C
C GENERAL WORK STORES.
C
C COMMON IWORK1,IWORK2,WORK3,WORK4,WORK5,WORK6,WORK7,WORK8,WORK9,
1 WORK10,WORK11,WORK12,WORK13,WORK14,WORK15,WORK16,WORK17,
2 WORK18,WORK19,ISCOLL
C
C DIMENSIONS ETC.
C
C COMMON HITE,FRAD,BRAD,ANGLE,TANGLE,DIST,STRTE,CNPHI(64),FLTPTH
C
C CONTROL PARAMETERS.
C
C COMMON IRECRG,IRGRA,IRGRB,JOBFIN,NOSAMP,SAMPLE,ANEUNO,ENSENCE,IFB4
C
C

```



```

C MAGGIE 3A 1566
C 20.7.64 A.D.PURNELL 1570
C 9.68 MODIFIED NBS VERSION FOR FORTRAN V. A.C.B.RICHARDSON NBS1570A
C FOR TRUNCATED CONE SCATTERING SAMPLE, INCLUDING AUTOMATIC ELASTIC NBS1571
C SCATTERING CALCULATION AND GRAPHS. NBS1571A

INCLUDE CINC1
DIMENSION LAWTP(100)
ITER=1 1653
INIT = 16719 1654
READ(5,51) OCT
CALL RANDOM(OCT)
51 FORMAT(0I12)
READ(5,200)(HEAD(I),I=1,12)
WRITE(6,200)(HEAD(I),I=1,12)
200 FORMAT(12A6)
CALL DAT IN
CALL INPUT
WRITE(6,52) OCT
52 FORMAT(5SH0THE STARTING VALUE FOR THE RANDOM NUMBER GENERATOR IS
1012) 1656
103 NOSAMP=0
IRECRG=U
IRGRA=0
IRGRB=0
ANEUNO=0.0
THERM=0.25000001E-7
1 IF (IRGRA)2,7,4
2 PRINT 3,ANEUNO
3 FORMAT(1H1,4THERROR CONDITION-BIRTH REGISTER NEGATIVE ANEUNO=F6.0
2) 1658
IRGRA=0
IRGRB=0
GO TO 7
1669
IRGRA=IRGRA-11
1670
IRGRB=IRGRB-3
1671
1672

```

```

IWORK1=IBIRTH(IRGRB+1) 1673
MAT=IBIRTH(IRGRB+2) 1674
IWORK2=IBIRTH(IRGRB+3) 1675
EIN=BIRTH(IRGRA+1) 1676
WORK3=BIRTH(IRGRA+2) 1677
WORK4=BIRTH(IRGRA+3) 1678
WORK5=BIRTH(IRGRA+4) 1679
WORK6=BIRTH(IRGRA+5) 1680
WORK7=BIRTH(IRGRA+6) 1681
WORK8=BIRTH(IRGRA+7) 1682
WORK19=BIRTH(IRGRA+8) 1683
UIN=BIRTH(IRGRA+9) 1684
WIN=BIRTH(IRGRA+10) 1685
WORK11=BIRTH(IRGRA+11) 1686
IF(IRGRA-550)13,5,13 1687
IF(IRECRG)13,13,6 1688
IRECRG=IRECRG-11 1689
LOC=LOC-5600 NBS1690
READ(37)IRECRG,(BIRTH(I),I=551,4950),(IBIRTH(I),I=151,1350) NBS1692
IRGRA=IRGRA+4400 1693
IRGRB=IRGRB+1200 1694
GO TO 13 1695
IF(ANEUNO-SAMPLE)10,8,8 1696
NOSAMP=NOSAMP+1 1697
ANEUNO=0.0 1698
CALL OUTPUT 1699
IF(ITEMG-1)81,82,81 1700
81 IF(NOSAMP-JOBFIN)10,9,9 1701
82 IF(NAC1(1)-2)9,84,9 1702
84 IF(NOCOU-33)9,85,9 1703
85 CALL AUTO 1704
IF(ITER-10) 103,9,9 1705
9 CONTINUE 1706
CALL END 1715
10 CALL CRNEU 1716
IF(INSENSE)11,13,11 1717

```

```

11 PRINT 12 FORMAT(120H0SERIAL STEP EN LETH WT L
12   X Y Z TIME FATE ISO. ACT. RAN. LAW N
13   IF (NSENSE) 14,15,14 N NBS1719
14   PRINT 2,WORK6,WORK7,WORK8,WORK19 1720
15   CALL TRACK 1721
16   WORK1=WIN 1722
17   ISCOLM=ISCOLL 1723
18   GO TO (17,16,22),ISCOLM 1724
19   PESCE=PESC+WORK1 1725
20   GO TO 18 1726
21   CESC=CESC+WORK1 1727
22   IF (NSENSE) 19,20,19 1728
23   PRINT 34,ANEUMO,IWORK1,EIN,UIN,WIN,WORK3,WORK4,WORK
24   25,WORK16,WORK17,WORK18,WORK19,ISCOLM 1729
25   CONTINUE 1730
26   505 IF (IWORK1) 506,506,507 1731
27   506 ANOCOL=ANOCOL+WORK1 1732
28   GO TO 1 1733
29   507 IF (IWORK1-1) 1,508,509 1734
30   508 I=ICOLL(1) 1735
31   ONCOL(I)=ONCOL(I)+WORK1 1736
32   GO TO 1 1737
33   509 NOIN=0 1738
34   DO 511 I=1,1,WORK1 1739
35   IF (ICOLL(I)-2) 511,511,510 1740
36   NOIN=NOIN+1 1741
37   LAWTP(NOIN)=LAWTP(I) 1742
38   CONTINUE 1743
39   IF (NOIN=2) 518,512,518 1744
40   IF (LAWTP(I)-LAWTP(2)) 513,514,514 1745
41   I=LAWTP(2) 1746
42   J=LAWTP(1) 1747
43   GO TO 515 1748
44   I=LAWTP(1) 1749

```

```

J=LAWTP(2) 1763
515 IF(I-5)516,516,518 1764
516 IF(J-5)517,517,518 1765
517 TOCOL(I,J)=TOCOL(I,J)+WCRK1 1766
518 IF(ICOLL(I)-2)524,519,519 1767
519 J=ICOLL(I)-1 1768
NOIN=0 1769
DO 521 I=1,1WORK1 1770
IF(ICOLL(I)-2)521,520,520 1771
NOIN=NOIN+1 1772
520 CONTINUE 1773
IF(NOIN-1)522,522,523 1774
521 DEGENE(J)=DEGENE(J)+WORK1 1775
NO TO 1 1776
DEGENI(J)=DEGENI(J)+WORK1 1777
60 TO 1 1778
NBS1778
1779
DO 526 I=1,1WORK1 1780
IF(ICOLL(I)-2)526,525,525 1781
NOIN=NOIN+1 1782
NOIN1=ICOLL(I) 1783
522 CONTINUE 1784
IF(NOIN-1)527,529,528 1785
523 ALLEL=ALLEL+WORK1 1786
60 TO 1 1787
NBS1787
1788
DO 528 I=1,1WORK1 1789
IF(ICOLL(I)-2)528,527,529 1790
EDGEN(J)=EDGEN(J)+WORK1 1791
60 TO 1 1792
NBS1792
1793
1794
1795
1796
1797
1798
1799
C COLSN
22 1WORK1=1WORK1+1
CALL CR
IF(IFB4)23,24,23

```

```

23      CALL CALC
24      CALL NAPAN
25      WORK1=UOUT
26      IF(IWORK1-100)25,25,26
27      IF(IREACT(1)-2)532,531,532
28      ICOLL(IWORK1)=1
29      IF(IWORK1-1)26,5310,26
30      ALAM(1)=ALAM(1)+WORK1
31      GO TO 26
32      IF(NREACT(1)-16)534,533,533
33      ICOLL(IWORK1)=2
34      IF(IWORK1-1)26,5330,26
35      ALAM(2)=ALAM(2)+WORK1
36      GO TO 26
37      DO 535 1=1,NLWREV
38      IF(ENOLW(1)-LAWREV(1))535,536,535
39      CONTINUE
40      ICOLL(IWORK1)=3
41      LAWTP(IWORK1)=1
42      IF(IWORK1-1)26,5350,26
43      ALAM(3)=ALAM(3)+WORK1
44      GO TO 26
45      ICOLL(IWORK1)=ELWNO(1)+3
46      LAWTP(IWORK1)=ELWNO(1)+1
47      IF(IWORK1-1)26,5360,26
48      ALAM(1+3)=ALAM(1+3)+WORK1
49      IF(NSECS)32,507,27
50      WORK10=UOUT
51      DO 31 1=1,NSECS
52      WORK12=LOUT(1)
53      WORK9=UOUT(1)
54      IF(WORK12-THERM)28,29,29
55      WORK12=THERM
56      WORK9=LOG(THERM/0.25E-7)
57      ANGLE=COSEPH(1)
58      CALL TWIST(ANGLE)
59

```

```
1836 CALL WR1BT
1837 IF (NSENSE) 30, 31, 30
1838 PRINT
1839 2WORK14, WORK15, WORK16, WORK17, WORK18, WORK19, WORK13,
1840 3, NORNGE(I), NOLAW(I), NSECS
1841 CONTINUE
1842 GO TO 1
1843 FORMAT(1H0, F7.0, I4, 9F7.3, F10.4, I5, 3I6, I5, I7)
1844 FORMAT(1X, F7.0, I4, 9F7.3, F10.4, I5, 3I6, I5, I7)
1845 END
```

SUBROUTINE INPUT

```

0251
0252
0253
0254
0255
0256
0257
0258
0259
0260
0261
0262
0263
0264
0265
0266
0267
0268
0269
0270
0271
0272
0273
0274
0275
0276
0277
0278
0279
0280
0281
0282
0283
0284
0285
0286

0251 READ(5,201)NSTDP
0252 READ(5,202)(ANGS(I),STDIST(I),I=1,NSTDP)
0253 WRITE(6,2010)(ANGS(I),STDIST(I),I=1,NSTDP)
0254 2010 FORMAT(31HUSTART DISTRIBUTION PROVIDED IS /9(6F10.4/))
0255 2011 FORMAT(6F10.4)
0256 2012 FORMAT(6F10.4)
0257 C DEFINE START DISTRIBUTION TO INCLUDE THE MINIMUM ANGLE WHICH WILL
0258 JUST ENCLOSE THE SAMPLE (USUALLY ABOUT 10-16 DEGREES) BY UP TO 25 NBS0258
0259 POINTS.
0260 DO 211 I=1,NSTDP
0261 ANGS(I)=COS(ANGS(I)/57.295828)
0262 CONTINUE
0263 DO 212 I=1,NSTDP
0264 IF(THETAM-ANGS(I))212,213,213
0265 CONTINUE
0266 WRITE(6,2120)
0267 FORMAT(41HINSUFFICIENT START DISTRIBUTION PROVIDED.)
0268 CALL EXIT
0269 AREA(1)=0.0
0270 STDIST(I)=STDIST(I-1)-(STDIST(I-1)-STDIST(I))*(ANGS(I-1)-THETAM)/(
0271 2ANGS(I-1)-ANGS(I))
0272 ANGS(I)=THETAM
0273 L=I-1
0274 LI=I-1
0275 DO 214 K=1,L
0276 AREA(K+1)=AREA(K)+((STDIST(K+1)+STDIST(K))/2.0)*(ANGS(K)-ANGS(K+1))
0277 2)
0278 CONTINUE
0279 DO 215 K=1,64
0280 AIEK
0281 SEG=(AREA(I)*(2.0*AI-1.0))/128.0
0282 DO 216 L=1,I
0283 IF(CAREA(L)-SEG)216,217,218
0284 CONTINUE
0285 WRITE(6,2160)
0286 2160 FORMAT(63HERROR IN INPUT. PARTIAL SUM LRSS THAN THE TOTAL. CALL

```



```

4020 LGR=M
  IF(S-EMC(M))4030,50000,4040
4030 M=M-32
  GO TO 4050
4040 M=M+32
4050 LGR=M
  IF(S-EMC(M))4060,50000,4070
4060 N=M-16(1)
  GO TO 4080
4070 M=M+16
4080 LGR=M
  IF(S-EMC(M))4090,50000,4100
4090 M=M-8
  GO TO 4110
4100 M=M+8
4110 LGR=M
  IF(S-EMC(M))4120,50000,4130
4120 M=M-4
  GO TO 4140
4130 M=M+4
4140 LGR=M
  IF(S-EMC(M))4150,50000,4160
4150 M=M-2
  GO TO 4170
4160 M=M+2
4170 LGR=M
  IF(S-EMC(M))4180,50000,4190
4180 M=M-1
  GO TO 4200
4190 M=M+1
4200 LGR=M
  IF(S-EMC(M))4210,50000,5000
5000 M2=1DATA(M1)+LGR-MISSM
  AVPATH=DATA(N2)
  FLUXFT=FLUXFT/AVPATH

```

```

AVCOS=(CNPH1(32)+CNPH1(33))/2.0
DO 2150 I=1,LI
  IF (ANGS (I+1)-AVCOS) 2151,2152,2150
  CONTINUE
  I=LI
2151 FACT=STDIST(I)-(STDIST(I+1))*(ANGS (I)-AVCOS)/(ANGS (I)
2-ANGS (I+1))
  GO TO 2153
2152 FACT=STDIST(I)
2153 FLUXFT=FLUXFT*STDIST(I)/FACT
  READ(5,201)NLWREV
  READ(5,220)(LAWREV(I),I=1,NLWREV)
  WRITE(6,2193)
2193 FORMAT(38HOLAW REFERENCE NUMBER M.C.LAW NUMBER.)
  WRITE(6,2194)(LAWREV(I),I=1,NLWREV)
2194 FORMAT(10,10X,110)
  IF(LFB4)2192,234,2192
2192 READ(5,201)IANALA,IANALB,IANALC,IANALD
  READ(5,220)
  READ(5,220)(NUCL(I),I=1,NANAL)
220 FORMAT(2110)
  WRITE(6,221)ANAL
221 FORMAT(15H054 ANALYSIS ON I3,10H ACTIONS.)
  WRITE(6,222)(NUCL(I),I=1,NANAL)
222 FORMAT(15HNUCLIDE NOS. 24I4)
  WRITE(6,223)(IANAL(I),I=1,NANAL)
223 FORMAT(16HOK.P.ACTION NOS.,I3,23I4)
  M3=I-NMATS
  IXNUC1=IDATA(M3)
  M4=M3+4*NMATS
  NUCS=IDATA(M4)
  DO 3 I=1,IANAL
  M5=IXNUC1+NUCL(I)
  M6=8*NUCS+M5
  NOACA(I)=IDATA(M6)
  IXACT1(I)=IDATA(M5)
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K=1
NOACA1=NOACA(1)+IXACT1(I)
DO 4 J=1,NOACA1
L=MMM+K
IF (LDATA(L)-NAC1(I))7,5,7
K=K+1
CONTINUE
4      WRITE(6,6)
      CALL EXIT
      NACT(I)=K
CONTINUE
5      FORMAT(18H K.P. ACTION NUMBER,14,29H DOES NOT APPEAR IN THE DATA.)
      WRITE(6,224) (NACT(I),I=1,NANAL)
224    FORMAT(15HUM.C.ACTION NO.,'24I4)
      READ(5,202)(FCCOS(I),I=1,33)
      33 COSINES WITH 33 VALUES OF THE ANGULAR DISTRIBUTION MUST BE
      SUPPLIED, WITH MONOTONICALLY INCREASING VALUES OF THE COSINE.
      DO 225 J=1,NANAL
      READ(5,202)(FCVAL(J,I),I=1,33)
      AREB=0.0
      DO 228 I=1,32
      AREB=AREB+ABS (FCCOS(I+1)-FCCOS(I))*(FCVAL(J,I)+FCVAL(J,I+1))/2.0NBS0381
2)    CONTINUE
      1B=15
      IA=1
      1B=15
      WRITE(6,226)(FCCOS(I),I=IA,IB)
226    FORMAT(15HOCOSINES ,15F7.3)
      WRITE(6,227)(FCVAL(J,I),I=IA,IB)
227    FORMAT(15H DISTRIBUTION ,15F7.3)
      DO 229 I=IA,IB
      FCVAL(J,I)=FCVAL(J,I)/AREB
      CONTINUE
      WRITE(6,230)(FCVAL(J,I),I=IA,IB)
229    FORMAT(15H UIFF. X SECT. ,15F7.3)
230

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0395 IF (IA-1)2300,2300,2301
2300 IA=16 0396
     IB=30 0397
     GO TO 2<80 0398
2301 IF (IA-16)2302,2302,225 0399
2302 IA=31 0400
     IB=33 0401
     GO TO 2280 0402
225  CONTINUE 0403
     IMINM=1 0404
     DO 252 L=1,IMINM 0405
     IMIN(L)=IMINM
     READ(5,250)IMAX(L) 0406
250  FORMAT(1I10) 0407
     PRINT 251,IMAX(L),L 0408
251  FORMAT(1UHUTHERE ARE,12,46H SUPPLEMENTARY RANGES ASSOCIATED WITH ANBS0410
2NALYSIS,12) 0411
     IF (IMAX(L)>252,252,253 0412
     IMAXM=IMAX(L)-1+IMINM 0413
253  DO 254 I=IMINM,IMAXM 0414
     READ(5,202)ENVAL(L,I) 0415
     READ(5,202)(SUPVAL(I,K),K=1,33) 0416
     AREB=0.0
     DO 258 J=1,32 0417
     AREB=AREB+ABS (FCCOS(J+1)-FCCOS(J))*((SUPVAL(I,J)+SUPVAL(I,J+1))/2NBS0419
2.0)
258  CONTINUE 0418
     IA=1 0420
     IB=15 0421
     259  WRITE(6,258)ENVAL(L,I) 0422
2580 FORMAT(25H0 THIS RANGE APPLIES BELOW,F7.4,5H MEV.) 0423
     WRITE(6,226)(FCOS(J),J=IA,IB) 0424
     WRITE(6,227)(SUPVAL(I,J),J=IA,IB) 0425
     DO 260 J=IA,IB 0426
     SUPVAL(I,J)=SUPVAL(I,J)/AREB 0427
     CONTINUE 0428
260

```

```

      WRITE(6,230)(SUPVAL(I,J),J=1A,1B) 0431
      1F((IA-1)261,261,262 0432
261    IA=16 0433
          IB=30 0434
          GO TO 259 0435
262    IF((IA-16)263,263,254 0436
263    IA=31 0437
          IB=33 0438
          GO TO 259 0439
254    CONTINUE 0440
          IMIN=IMAX+1
          IF((IMAX-24)252,252,256 0441
          WRITE(6,257) 0442
          FORMAT(27H1IMAX HAS EXCEEDED 24-HALT.) 0443
          CALL EXIT 0444
252    CONTINUE 0445
          READ(5,201)NOCOU 0446
          READ(5,202) (COUNT(I),I=1,NOCOU) 0447
          WRITE(6,232) NOCOU,(COUNT(I),I=1,NOCOU)
232    FORMAT(10H0THERE ARE,13,56H COUNTERS LOCATED AROUND THE EQUATOR OFNBS0450
          2 THE SCATTERER AT/(10F9.4) 0449
          READ(5,201) NENSP 0451
          READ(5,202) (ENSP(I),I=1,NENSP) 0452
          WRITE(6,233) (ENSP(I),I=1,NENSP) 0453
233    FORMAT(48H0THE ENERGY SPECTRA OF THE B4 ANALYSIS CELLS ARE/(10F9.3
          2) 0454
234    DO 2340 K=1,32 0455
          FCCOSD(K)=FCCOS(K+1)-FCCOS(K) 0456
2340  CONTINUE 0457
          RETURN 0458
END 0459
          0460
          0463

```

C TRUNCATED CONE SCATTERING SAMPLE, INCIDENT NEUTRONS ALONG THE AXISNBS

```

INCLUDE CINC1
DIMENSION TEMP(10),TRIG(4)
61  TEMPENRANDA(64)
WORK3=CEVPHI(ITEMP)
62  TEMP(1)=ERANDA(-1)
TEMP(2)=ERANDA(-2)
TEMP(3)=TEMP(1)*TEMP(1)
TEMP(4)=TEMP(2)*TEMP(2)
TEMP(5)=TEMP(3)+TEMP(4)
IF(TEMP(5)>62)62,62,63
IF(TEMP(5)-1.0)64,62,62
63  TRIG(1)=(TEMP(3)-TEMP(4))/TEMP(5)
TRIG(2)=(2.0*TEMP(1)*TEMP(2))/TEMP(5)
TRIG(3)=SQRT((1.0-WORK3*WORK3)
WORK4=TRIG(1)*TRIG(3)
WORK5=TRIG(2)*TRIG(3)
64
C DOES THE NEUTRON HIT THE TARGET
C
TEMP(1)=FRAU+(DIST+HITE)*TRIG(3)/WORK3
IF(TEMP(1)>67.67,67.65
65  TEMP(2)=(DIST+HITE)/WORK3
66  WORK6=-HITE
WORK7=WORK4*TEMP(2)
WORK8=WORK5*TEMP(2)
ANEUNO=ANEUNO+1.0
EIN=STRTE
UIN=LOG(EIN/0.25E-7)
MAT=1
IWORK1=0
WORK19=0.0
WIN=1.0
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NBS3121
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EU.3124
EU.3126
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NBS3150
NBS3151
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WORK11=0.0
WORK2=0
RETURN
AMISS=AMISS+1.0
GO TO 61
END

67

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SUBROUTINE TRACK

C TRACKS NEUTRONS IN A TRUNCATED CONE WITH AXIS LYING ALONG THE NEUTRON BEAM

```

INCLUDE CINC1
CALL E6MV
PATHL=(-(LOG(RANDA(1)))*PATH
IF (WORK3)1,6,2
1 TRS=-(HITE+WORK6)/WORK3
RAD2=FRAD**2
GO TO 3
2 TRS=(HITE-WORK6)/WORK3
RAD2=BRAD**2
3 WORK17=WORK7+TRS*WORK4
WORK18=WORK8+TRS*WORK5
TRR2=WORK17**2+WORK18**2
IF (RAD2.LT.TRR2) GO TO 6
IF (TRS.GT.PATHL) GO TO 13
ISCOLL=2
IF (WORK3.GT.0) GO TO 4
WORK16=-HITE
GO TO 5
4 WORK16=HITE
5 WORK19=WORK19+TRS/SPEED
RETURN
6 CALL FONEC(PATHL)
IF (HITE.LE.ABS(WORK16)) GO TO 7
TRA2=WORK17**2+WORK18**2
RAD2=(FRAD+(HITE+WORK16)*TANGLE)**2
IF (RAD2.GT.TRA2) GO TO 14
7 ISCOLL=1
NRTS=0
RAD=FRAD+(HITE+WORK6)*TANGLE
TA=1.-WORK3**2*(1.+TANGLE**2)
TB=WORK7*WORK4+WORK8*WORK5-WORK3*TANGLE*RAD

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```

TC=WORK7**2+WORK8**2-RAD**2
IF(ABS(TA).GT.1.E-30)GO TO 8
IF(ABS(TB).LT.1.E-30)GO TO 21
TR=-TC/(2.*TB)
GO TO 12
8  TD=TB*TB-TA*TC
   IF(TD)15,10,9
9   TE=-TB+SQRT(TD)
   IF(TE.LT.0)GO TO 10
   NRTS=1
   TR=TE/TA
10  TE=-TB-SQRT(TD)
   IF(TE.LT.0)GO TO 11
   NRTS=NRTS+1
   TR=TE/TA
11  IF(NRTS-1)17,12,19
12  CALL F01EC(TR)
   WORK19=WORK19+TR/SPEED
   RETURN

13  CALL F01EC(PATHL)
14  ISCOLL=3
   WORK19=WORK19+PATHL/SPEED
   RETURN

C   ERROR EXITS

15  WRITE(6,16)
16  FORMAT(24H IMAGINARY ROOT IN TRACK)
   GO TO 23
17  WRITE(6,18)
18  FORMAT(26H NO POSITIVE ROOT IN TRACK)
   GO TO 23
19  WRITE(6,20)
20  FORMAT(28H TWO POSITIVE ROOTS IN TRACK)
   GO TO 23

```

```
21 WRITE(6,22)
22 FORMAT(22H INDETERMINATE ROOT IN TRACK)
23 WRITE(6,24) WORK3,WORK4,WORK5,WORK6,WORK7,WORK8,PATHL
24 FORMAT(7F10.3)
      CALL EXIT
      END
```

```

INCLUDE CINC1
DIMENSION FACT(33), PATHLL(33), ANGLEE(33), KL(33)
NENSP1=NENSP-1
M1=IXMAT+NMAT+NMAT
NEM1+4*NMATS
MISSM=IDATA(1)
K2=IDATA(M1)-MISSM
DO 1 I=1, NOCOU
  COSOMECOUNT(I)
  CALL FPATH(COSOM, PATHL, ANGLEF)
  DO 2 K=1, 32
    IF (FCCOS(K+1)-ANGLEF) 2, 3, 3
    CONTINUE
 2
  K=32
  AFCOS=FCCOS(K)
  FACT(I)=(ANGLEF-AFCOS)/FCCOSD(K)
  PATHLL(I)=PATHL
  ANGLEE(I)=ANGLEF
  KL(I)=K
  CONTINUE
  DO 7 J=1, NATOM
    NFORM=0
    NAC=NAC1(J)
    NUCLID=NUCL(J)
    PART=PART1(J)
    ATOM=ATOM1(J)
    WSQ=ATOM*ATOM-1.0
    DO 6 I=1, NOCOU
      KEKL(I)
      AS=ANGLEE(I)
      FACTOR=FACT(I)
      PATHL=PATHLL(I)
      ASQ=AS*AS
      IF (PART) 7, 7, 8
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4200 LGR=M
  IF (S-EMC(M)) 4210,5000,5000
4210 LGR=N-1
5000 M3=M2+LGR
  PATHB=(-DATA(M3))
  QU=J(J)
  PTH=PATHL/PATHB
  ADD=QU*SIGMA*EXP (PTH)
  DEA(I,L)=DEA(I,L)+ADD
  IF (NAC.NE.2) GO TO 24
  IF (IWORK1.LE.1) GO TO 23
  I1=IWORK1-1
  DO 22 I1=1,I1
  IF (ICOLL(I1).NE.1) GO TO 24
  22 CONTINUE
  23 DEB(I,L)=DEB(I,L)+ADD
  IF (IWORK1.LE.1) GO TO 0
  DED(I,L)=DED(I,L)+ADD
  GO TO 6
  24 DEC(I,L)=DEC(I,L)+ADD
  6  CONTINUE
  7  CONTINUE
  GO TO 40
  38 PRINT 39,REDEN
  39 FORMAT(1H0,F10.4,32hNEGATIVE ENERGY COMPUTED IN CREN)
  CALL EXIT
  40 RETURN
  END

```

SUBROUTINE FPATH(COSOM,PATHL,ANGLEF)

C PATHLENGTH AND ANGLE FOR B4 ANALYSIS. FOR TRUNCATED CONE. SAMPLE.
C AXIS ALONG THE INCIDENT NEUTRON BEAM.

INCLUDE CINC1

C NEUTRON OUT AT 90 DEGREES.
IF(ABS(COSOM).GE.1.E-5)GO TO 1
ANGLEF=WORK4
TSQ=ANGLE**2
XPR=FRAU/TANGLE+HITE+WORK16
PATHL=SQRT(XPR*XPR*TSQ-WORK18**2)-WORK17
RETURN

C NEUTRON OUT ENTRANCE FACE
1 SINOMESQRT(1.-COSOM**2)
ANGLEF=COSOM*WORK3+SINOM*WORK4
IF(COSOM.GT.0)GO TO 2
IF(FRAD.LT.ABS(WORK18))GO TO 3
FTAN=(SQRT(FRAD**2-WORK18**2)-WORK17)/(HITE+WORK16)
TANOM=SINOM/COSOM
IF(FTAN+TANOM.LT.0)GO TO 3
PATHL=(HITE+WORK16)/COSOM
RETURN

C NEUTRON OUT EXIT FACE
2 BTAN=(SQRT(BRAD**2-WORK18**2)-WORK17)/(HITE-WORK16)
TANOM=SINOM/COSOM
IF(BTAN.LT.TANOM)GO TO 3
PATHL=(HITE+WORK16)/COSOM
RETURN

C NEUTRON OUT THE SIDE. BUT NOT AT 90 DEGREES. NORMAL SOLUTION.
3 TSQ=ANGLE**2
XPR=FRAU/TANGLE+HITE+WORK16

```
PA=SINOM**2-(COSOM**2)*TSQ
PB=WORK17*SINOM-XPR*COSOM*TSQ
PC=WORK17**2-(XPR**2)*TSQ+WORK18**2
IF (ABS(PA)*LT.1.E-30) GO TO 4
PA=THL=(SQR1(PB**2-PA*PC)-PB)/PA
RETURN
```

```
C   LINEAR SOLUTION
4  PATHL=-PC/(2.*P13)
RETURN
END
```

```

C          SUBROUTINE FPATH(COSOM,PATHL,ANGLEF)
C          PATH LENGTH AND ANGLE FOR THE B4 ANALYSIS, FOR A CYLINDRICAL SAMPLE
C          WITH ITS AXIS ALONG THE INCIDENT NEUTRON BEAM.

INCLUDE CINC1
PATHL=SQRT(FRAD**2-WORK18**2)-WORK17
IF(ABS(COSOM).GT.1E-5)GO TO 1
ANGLEF=WORK4
GO TO 3
1 SINOM=SQRT(1-COSOM**2)
TANOM=SINOM/COSOM
ANGLEF=COSOM*WORK3+SINOM*WORK4
IF(COSOM.GT.1E-5)GO TO 2
FTAN=PATHL/(HITE+WORK16)
IF((FTAN+TANOM).LT.1E-5)GO TO 3
PATHL=(HITE+WORK16)/COSOM
RETURN
2 FTAN=PATHL/(HITE-WORK16)
IF(FTAN.LT.1ANOM)GO TO 3
PATHL=(HITE+WORK16)/COSOM
RETURN
3 PATHL=PATHL/SINOM
RETURN
END

```

C SUBROUTINE OUTPUTS RESULTS AT THE END OF EACH SAMPLE.

```

INCLUDE CINC1
DIMENSION P(33)
PRINT 1, (HEAD(I), I=1,12)
1 FORMAT(12A6)          2, NO SAMP
2 FORMAT(1H0,52X,13HSAMPLE NUMBER,I3)      3, SAMPLE
3 FORMAT(1H0,40X,F8.0,33H NEUTRONS STARTED IN EACH SAMPLE.)
PART = NO SAMP
TOT = PART * SAMPLE
PRINT 5, TOT
5 FORMAT(1H0,39X,F9.0,35H NEUTRONS HAVE BEEN TRACKED SO FAR.)
PRINT 4, PESC
4 FORMAT(24HUPLANE SURFACE ESCAPES =,F11.3/24H0CURVED SURFACE ESCAPE
25=F11.3)
CALL OUTB3
II = NENSP - 1
IF (IFB4 .EQ. 0) GO TO 7
CALL OUTB4 (DEA,1,IANALA)
IF (IANALA .LE. 0) GO TO 8
DO 6 II = 1,II
P(I) = DEA (IANALA,I)
CALL APLOT3 (P,II,ENSP,IANALA)
CALL OUTB4 (DEB,2,IANALB)
IF (IANALB .LE. 0) GO TO 10
DO 9 II = 1,II
P(I) = DEB (IANALB,I)
CALL APLOT3 (P,II,ENSP,IANALB)
CALL OUTB4 (DEC,3,IANALC)
IF (IANALC .LE. 0) GO TO 12
DO 11 II = 1,II
P(I) = DEC (IANALC,I)
CALL APLOT3 (P,II,ENSP,IANALC)
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NBS 3994
NBS 3996

```

```
12 CALL OUT34(JED,4,IANALD)
  IF(IANALD.LE.0) GO TO 7
  DO 13 I=1,I1
  P(I)=JED(IANALD,I)
  CALL APLOT3(P,I1,ENSP,IANALD)
  RETURN
13 END
```

13

7

SUBROUTINE AUTO

C TO ALLOW AUTOMATIC ITERATION UNDER THE FOLLOWING CONDITIONS. 1978
 C NUMBER OF COUNTERS AND INPUT POINTS = 33. 1979
 C ONE NUCLIDE ONLY. 1980
 C ELASTIC CORRECTION ONLY. 1981
 C PHYSICAL ITERATION.
 C $(\Delta\sigma)/(\sigma)$ IS THE CHANGE IN THE CROSS SECTION EXCLUDING
 C MULTIPLE SCATTERING.

```

INCLUDE CINCL
DIMENSION TOTEL(33),CHI(33),ELM2(33),XIP3(33)
SUMEL=(DEB(1,33)+DEB(33,33))/32.0
DO 4 J=2,32
SUMEL=SUMEL+IEE(J,33)/16.0
4 CONTINUE
IF (ITER.NE.1) GO TO 300
READ(5,101)INTERM,IFCONT
101 FORMAT(2I10)
INTERM=INTERM-1
IF (IFCONT.EQ.0) GO TO 100
READ(5,104)(EZP(I),I=1,33)
104 FORMAT(6F10.4)
SUMEZP=(EZP(1)+EZP(33))/32
DO 107 J=2,32
107 SUMEZP=SUMEZP+EZP(J)/16
DO 108 J=1,33
108 EZP(J)=EZP(J)/SUMEZP
GO TO 102
100 DO 105 I=1,33
105 EZP(I)=FCVAL(1,I)
105 CONTINUE
102 IF (INTERM.LE.1) GO TO 300
106 READ(5,103)(SAMP(I),I=1,INTERM)
103 FORMAT(6F10.2)
300 CHISQ=0.0
      
```

```

DO 303 J=1,33
TOTEL(J)=DEB(J,33)/SUMEL
CHI(J)=(TOTEL(J)-EZP(J))*2)/EZP(J)
CHISQ=CHISQ+CHI(J)
DEDNJ=DED(J,33)/SUMEL
ELM2(J)=FCVAL(1,J)
208 NBS
2088 NBS2089
2090 NBS
2097 NBS
FACTORE=(TOTEL(J)-DEDNJ)/ELM2(J)
IF(ITER.GT.2) GO TO 302
XIP3(J)=(EZP(J)-DENNJ)/FACTOR
GO TO 303
302 XIP3(J)=FCVAL(1,J)/FACTOR
303 CONTINUE
2098 NBS
2099 NBS
2100 NBS
2101 NBS
2102 NBS
2103 NBS
2104 NBS
2105 NBS
214 SUM=(XIP3(1)+XIP3(33))/32.0
DO 207 J=2,32
SUM=SUM+XIP3(J)/16.0
207 CONTINUE
208 1=1,33
FCVAL(1,1)=(XIP3(1))/SUM
208 CONTINUE
209 1=1,33
DO 220 1=1,33
XIP3(1)=DEB(1,33)-DED(1,33)
SUMSEL=(XIP3(1)+XIP(33))/32
220 NBS
221 1=2,32
DO 221 1=2,32
SUMSEL=SUMSEL+XIP3(1)/16
221 NBS
222 1=1,33
DO 222 1=1,33
XIP3(1)=(XIP3(1)/SUMSEL-ELM2(1))/ELM2(1)
222 NBS
2106 NBS
2107 NBS
2108 NBS
2109 NBS
2110 NBS
2111 NBS2112
2112 NBS2113
2113 NBS2114
2114 NBS2115
2115 NBS
2116 ITER1=ITER+1
PRINT 210,ITER,ITER,ITER,ITER
210 FORMAT(1H0,7H COUNTER,4X,6HCOSINE,11X,3H EXP,7X,10H OUTPUT NO.,12,4X,
13H NO.,12,5H CHI,7X,3H NO.,12,11X,3H NO.,12,4X,11H DELTA SIGMA,2X,
27H COUNTER/2X,6H NUMBER,5X,5HANGLE,10X,5H INPUT,7X,11H ALL ELASTIC,6X,
37HSQUARED,8X,5H INPUT,11X,5H INPUT,9X,6H/SIGMA,2X,6H NUMBER)
PRINT 111,(J,COUNT(J),EZP(J),TOTEL(J),CHI(J),ELM2(J),FCVAL(1,J),

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NBS2116
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1XIP3(J),J=1,33)
111 FORMAT(14,7F15.6,14)
112 PRINT 112,CHISQ
112 FORMAT(1H0,3HCHISQ = 'F10.6)
997 IF(ITER-INTERM)415,416,416
415 SAMPLE=SAMP(ITER)
ITER=ITER+1
60 TO 999
416 ITER=10
60 TO 999
999 CALL DATIN
998 RETURN
END
```

SUBROUTINE RANDOM(ISTART)

```

DATA MASK/0377777777777777/
MULTA/30517578125/,C/.29103830E-10/
IRSTART
RETURN
ENTRY RDM(R)
IRE=AND(MASK,IR*MULTA)
K=C*(FLCAT(IR))
RETURN
END

```

FUNCTION RANDA(I)

```

CALL RDM(R)
RANDA=R
IF(I)1,1,3
1 RANDA=RANDA*2.-1.
3 RETURN
END

```

EU.4009

```

NBS4010
NBS4010A
EU.4011
EU.4012
EU.4013
EU.4014

```

FUNCTION NRANDA(I)

```

CALL RDM(XRAND)
IF(XRAND)10,10,1
1 IF(XRAND-1.)20,30,30
10 NRANDA=1
GO TO 40
20 NRANDA=XRAND*1+1
GO TO 40
30 NRANDA=1
40 RETURN
END

```

EU.0002

```

NBS0003
EU.0004
EU.0005
EU.0006
EU.0007
EU.0008
EU.0009
EU.0010
EU.0011
EU.0012

```

A Brief Description of the Components of MAGGIE

1. MAGGIE (Main Program). This program calls the various subroutines required for the analysis, retrieves neutrons from disc storage when required, outputs track parameters if desired, and records the various fates of tracked neutrons.
2. RANDOM (and RDM). This subroutine contains the random number generator, RDM, as a separate entry. RANDOM is called at the beginning of program MAGGIE to enter the starting value for RDM.
3. DATIN serves to set all of the output arrays to zero, and calls ABSYND.
4. ABSYND reads the required nuclear data from the MOULD tape and puts it in encoded sequential storage in the array DATA-IDATA for use during the Monte Carlo tracking.
5. INPUT reads and processes samples and experimental angular distribution data from card input and calculates most of the flux attenuation factor.
6. CRNEU creates random incident neutrons at the entrance face of the sample, in accordance with the input source distribution.
7. TRACK tracks neutrons in the sample, specifying coordinates of collision or escape.
8. EGMV. This subroutine computes the mean free path, velocity, and lethargy group number.
9. FONEC calculates coordinates at the end of a track from initial position, direction cosines, and track length.
10. CR. This subroutine, using random sampling of the information stored by ABSYND from the MOULD data tape, determines all of the parameters of a collision.
11. CALC determines some constants used in subroutine NAPAN.
12. NAPAN. This subroutine calculates and scores the probability of detection at each detector for each collision in the sample.
13. FPATH calculates the path length in the direction of each detector for each collision.
14. CREN is an abbreviated version of CR used by NAPAN that determines only the neutron energy.

15. TWIST chooses new direction cosines after a collision.
16. WRTBT stores any secondary neutrons produced, for recall at the end of the current tracking.
17. OUTPUT, OUTB3, OUTB4 and APLOT3. These subroutines print the results of the calculation.
18. AUTO performs the calculations required for iterative correction of the elastic angular distribution, prints the current output, and calls DATIN to begin the next iteration.
19. END, and its entries EXIT and EEXIT designate normal vs. error exit conditions.
20. TAPLAB returns the tape logical unit label.
21. SRFORT is a subroutine for skipping tape records.

Input requirements for MAGGIE-NBS

<u>FORMAT</u>	<u>VARIABLE</u>	<u>COMMENTS</u>
1. 012	OCT	Octal starting value for RDM.
2. 12A6	HEAD(I)	Arbitrary heading. Column one should be a 1.
3. I10	NMATS	The number of materials in the sample. This is always = 1.
4. I10, E10.4	IDATA(MAT5) DATA(MAT4)	Number of different nuclides in the material. Density of the material.
5. I10, E10.4	IDATA(K2) DATA(K4)	Nuclide reference number (i.e. position on the MOULD tape: see output of MOULD for this). Proportion of this nuclide in the material. This card is repeated for each nuclide in the material.
6. 2I10	IFB4 NSENSE	Positive for B4 output. Positive for track print.
7. 3F10.4	HITE FRAD ANGLE	Length of the samples, in cm. Entrance radius of same, in cm. Half-angle of same, in radians.
8. 2I10	JOBFIN ITMAG	Number of independant samples (= 1 if ITMAG > 0). Positive for automatic iteration.
9. F10.4	SAMPLE	Number of neutrons in each sample.
10. F10.4	STRTE	Starting energy in MeV. If the starting energy is set = 0 a fission spectrum is assumed.
11. F10.4	DIST	Distance from the center of the sample to the source (negative), in cm.
12. I10	NSTD _P	No. of start distribution points to be read in (≤ 25).
13. 6F10.4	ANGS(I)	Three pairs to a card, the angle and start distribution for that angle. The start distribution need not be normalized, and this card is repeated until NSTD _P pairs are read in.

14. I10

NLWREV

The number of Law Reference numbers for inelastic scatter, i.e. the number of Law Reference numbers to different angular distributions on the MOULD output (P.C.N.'s* 4-15 only) (NLWREV \leq 50).

15. 2I10

LAWREV(I)
LWNO(I)

One pair per card, the above Law Reference numbers and the monotonically increasing "Monte Carlo" law numbers allocated. Any inelastic law not given here will be printed in the B3 section results under Law zero.

NOTE: The following cards are not required if IFB4 \leq 0.

16. 4I10

IANALA
IANALB
IANALC
IANALD

Four markers, for the tables: A) Complete multiple scatter analysis. B) Elastic events only. C) Inelastic events only. D) Multiple elastic events only. A negative marker suppresses the table. A positive marker = N produces, in addition to the table, a graph for the Nth counter.

17. I10

NANAL

Number of actions (index I below) (i.e. different neutron processes, such as elastic, inelastic from the 1st excited state, etc.) to be processed by the subroutine NAPAN.

18. 6F10.4

NUCL(I)
NAC1(I)

The nuclide reference number (as in card 5), and the P.C.N.* for each action. There are NANAL such pairs.

19. 6F10.4

FCCOS(J)

Thirty-three values of cosine, including -1 and +1, monotonically increasing.

20. 6F10.4

FCVAL(I,J)

Thirty-three values of the angular distribution (index J) corresponding to the above values of cosines. Repeat for each action (index I) in the same order as in 18.

NOTE: The following three items are repeated, as a group, for each action.

21. I10

IMAX(I)

Number (index L) of supplementary ranges for this action (even if zero).

* The "P.C.N.'s (Particular Classification Numbers) are listed starting on p.8 of AWRE report no. O 70/63, "The Aldermaston Nuclear Data Library as at May, 1963", K. Parker.

NOTE: The following two items are repeated for each supplementary range. If, for any action, there is no supplementary range these items are omitted.

22.	F10.4	ENVAL(I, L)	Upper limit of the range in MeV. These must be listed in order of decreasing value.
23.	6F10.4	SUPVAL(L, J)	Thirty-three values (index J) of the distribution (range L) at the cosine values FCCOS(J).
24.	I10	NOCOU	Number of counters (≤ 33).
25.	6F10.4	COUNT(I)	Cosines of counter locations.
26.	I10	NENSP	Number of output energy points (≤ 33). The output will be classified into the bins formed by these points.
27.	6F10.6	ENSP(I)	Values of output energy points in MeV.
28.	2I10	NTERM IFCONT	Number of iterations (≤ 10). If this is a continuation of a previous run, for further iterations this should be non-zero (see Section II.c.6.f).
29.	6F10.4	SAMP(I)	Number of neutrons for each iteration, <u>except</u> the first.

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